

Search for New Phenomena in Dijet Angular Distributions in Proton-Proton Collisions at $\sqrt{s} = 8$ TeV Measured with the ATLAS Detector

G. Aad *et al.**

(ATLAS Collaboration)

(Received 3 April 2015; published 4 June 2015)

A search for new phenomena in LHC proton-proton collisions at a center-of-mass energy of $\sqrt{s} = 8$ TeV was performed with the ATLAS detector using an integrated luminosity of 17.3 fb^{-1} . The angular distributions are studied in events with at least two jets; the highest dijet mass observed is 5.5 TeV. All angular distributions are consistent with the predictions of the standard model. In a benchmark model of quark contact interactions, a compositeness scale below 8.1 TeV in a destructive interference scenario and 12.0 TeV in a constructive interference scenario is excluded at 95% C.L.; median expected limits are 8.9 TeV for the destructive interference scenario and 14.1 TeV for the constructive interference scenario.

DOI: 10.1103/PhysRevLett.114.221802

PACS numbers: 13.85.Rm, 12.60.Rc

The search for an internal structure of fermions and new forces that might govern that structure is a major goal of modern particle physics. The most powerful probes are scattering experiments with large momentum transfer. Collisions of protons at the Large Hadron Collider (LHC) resulting in two energetic jets of particles (dijets) provide the largest momentum transfer currently available, and therefore the deepest probe.

The angular distribution of jets relative to the beam axis in events with high dijet invariant mass (m_{jj}) provides stringent tests of perturbative quantum chromodynamics (QCD) as well as theories of new phenomena. QCD calculations predict that dijet production, dominated by t -channel gluon exchange in the kinematic region of proton-proton (pp) collisions at the LHC, features steeply falling m_{jj} distributions and angular distributions peaked at $|\cos(\theta^*)| = 1$, where θ^* is the polar scattering angle in the two-parton center-of-mass frame. New phenomena, such as strong gravity [1,2] or new interactions [3–6] typically predict angular distributions which are more isotropic.

Previous studies of dijet angular distributions, at the CERN SPS [7,8], the FNAL Tevatron [9,10], the CERN LHC at $\sqrt{s} = 7$ TeV [11–16], and by the CMS Collaboration at $\sqrt{s} = 8$ TeV [17], have reported results consistent with the standard model (SM). This Letter reports on studies of dijet angular distributions in pp collisions at $\sqrt{s} = 8$ TeV in data with an integrated luminosity of 17.3 fb^{-1} collected with the ATLAS detector in 2012.

A detailed description of the ATLAS detector is published elsewhere [18]. The detector is instrumented over almost the entire solid angle around the pp collision point, with layers of tracking detectors, calorimeters, and muon detectors.

The jets are measured using a calorimeter system composed of different detector types covering different regions in η [19] and depth. The electromagnetic calorimeter is composed of liquid-argon sampling calorimeters, using lead as an absorber, and is split into a barrel ($|\eta| < 1.475$) and two end caps ($1.375 < |\eta| < 3.2$). The hadronic calorimeter is divided into a barrel and two extended barrels ($|\eta| < 1.75$), and two endcaps ($1.5 < |\eta| < 3.2$). The barrel and extended barrels are sampling calorimeters with steel as absorber and scintillator tiles as the active medium, while the hadronic end caps are liquid-argon calorimeters with copper as the absorber. In the very forward regions ($3.1 < |\eta| < 4.9$) there are liquid argon calorimeters with copper and tungsten absorbers.

The data are selected using a trigger that requires a single high- p_T [19] jet above one of eight thresholds, ranging from 25 to 220 GeV. Because of the high rate of jets at lower p_T , only a fraction of the events from the lower seven thresholds are stored.

Individual jets are reconstructed using the anti- k_r jet clustering algorithm [20,21] with radius parameter $R = 0.6$. The inputs to this algorithm are clusters [22] of calorimeter cells with energy depositions significantly above the noise. Jet four-momenta are constructed by the vectorial addition of clusters of cells, treating each cluster as a four-momentum with zero mass. The jet four-momenta are then corrected to the jet energy scale [23] as a function of η and p_T for various effects, the largest of which are the hadronic shower response, detector material distribution, and pileup events [24]. This is done using a calibration scheme based on samples of simulated events and validated with test-beam [25] and collision data [22] studies.

*Full author list given at the end of the article.

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The rapidity of a jet is defined as $y = \frac{1}{2} \ln[(E + p_z)/(E - p_z)]$, where E is the jet energy and p_z is the momentum component along the beam axis [19]. The scattering angle between two jets can be expressed using the variable $\chi = e^{|y_1 - y_2|} = e^{2|y^*|}$, where y_1 and y_2 are the rapidities of the two jets, and $y^* = \frac{1}{2}(y_1 - y_2)$. The rapidity boost of the dijet system with respect to the center of mass of the colliding protons is calculated as $y_B = \frac{1}{2}(y_1 + y_2)$.

For each trigger, the event is required to have a jet with p_T sufficient to achieve a trigger efficiency greater than 99.5%. For the lowest (highest) threshold trigger, this corresponds to $p_T > p_T^{\min} = 47(333)$ GeV. Events are required to have at least two jets, each with $p_T > 50$ GeV; the dijet system, defined as the two jets with largest p_T , is required to have $|y^*| < 1.7$, $|y_B| < 1.1$, and $m_{jj} > 600$ GeV (where the m_{jj} requirement avoids the kinematic bias in the angular distributions introduced by the minimum p_T requirement).

The detector covers the angular range $|y^*| < 1.7$, corresponding to $\chi < 30$. This interval is divided into 11 bins, with boundaries at $\chi_n = e^{(0.3 \times n)}$ for $n = 0$ to 10, approximating the segmentation of the calorimeter in $\Delta\eta$. The data are further binned coarsely in m_{jj} with the expectation that low- m_{jj} bins are dominated by QCD processes and that signals associated with new physics would be found in higher dijet invariant mass bins. The bin edges are chosen to optimize the expected sensitivity to the model of contact interactions. The highest dijet mass observed is 5.5 TeV.

The SM predictions are estimated using the PYTHIA8 [26] v8.160 event generator with the AU2 [27] underlying-event tune and the CT10 [28] parton distribution functions (PDF). The simulated events are propagated through a detector simulation [29] that uses the GEANT4 [30] simulation package. Pileup conditions vary as a function of the instantaneous luminosity and are taken into account by overlaying simulated minimum-bias events generated with PYTHIA8 onto the hard-scattering process such that the observed distribution of the average number of interactions per bunch crossing is reproduced. The same reconstruction and event selection are applied to the simulated events and the data.

The PYTHIA8 calculations are primarily to leading order (LO) in QCD with simulation of higher-order contributions included in the shower modeling. Events generated by PYTHIA8 are reweighted using a correction factor calculated based on the ratio of the next-to-leading-order (NLO) cross-section calculation from NLOJET++ [31–33] v4.1.2 to the LO+shower calculation from PYTHIA8:

$$K(\chi, m_{jj}) = \frac{\sigma_{\text{NLO}}(\chi, m_{jj})^{\text{NLOJET++}}}{\sigma_{\text{LO+shower}}(\chi, m_{jj})^{\text{Pythia8}}}.$$

The K factors decrease with χ and thus modify the shape of the angular distributions; the impact ranges from a few percent at low m_{jj} to approximately 15% for the highest

m_{jj} region. Additional processes accounting for electro-weak (EW) effects not included in PYTHIA8 (virtual weak boson exchange and Sudakov-type logarithms) are included as EW corrections [34]. The effect is most pronounced at high m_{jj} and low χ , and the correction factors range from unity at low m_{jj} to 0.98–1.12 in the highest m_{jj} region. The EW corrections and the NLO K factors are applied as a function of χ and m_{jj} to the PYTHIA8 prediction.

Figure 1 shows the distributions of the data as a function of χ . The distribution in each m_{jj} region is normalized to unity, as the sensitivity to new phenomena is due to the angular distribution rather than normalization. The predicted SM distributions are also shown in Fig. 1 and describe the data well. The EW corrections substantially improve the agreement of the SM prediction with data at high m_{jj} , as can be appreciated from the comparison of the

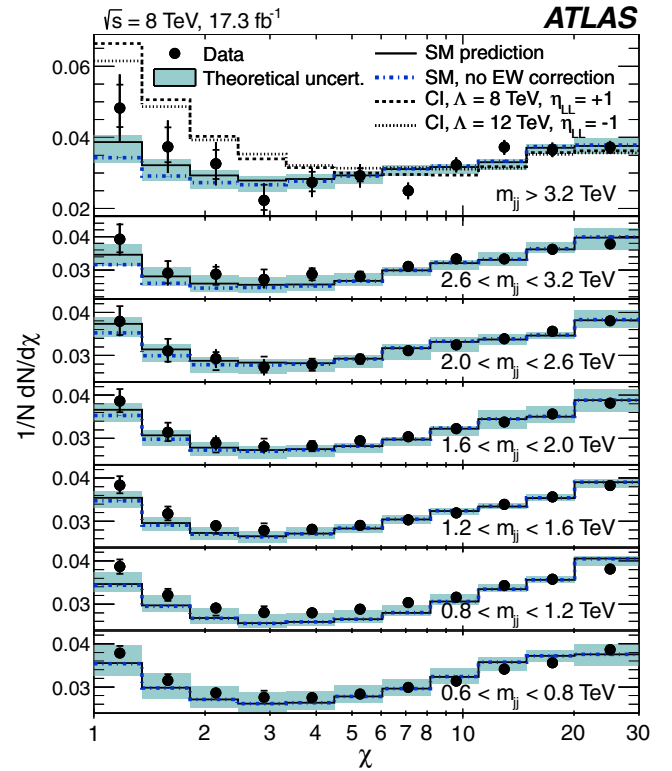


FIG. 1 (color online). Distributions of the dijet angular variable χ for several regions of dijet invariant mass m_{jj} . Shown are the data (dots), SM predictions (solid lines) with uncertainties, and predictions for a benchmark theory of new contact interactions with a left-chiral color-singlet coupling with constructive (dotted, $\eta_{LL} = +1$ and $\Lambda = 12$ TeV) or destructive (dashed, $\eta_{LL} = -1$ and $\Lambda = 8$ TeV) interference with the SM processes. Also shown is the impact of the electroweak (EW) corrections described in the text. The error bars on data represent the experimental and statistical uncertainties added in quadrature, with a tick representing experimental uncertainties only. The theoretical uncertainties are displayed as a shaded band around the SM prediction.

predictions with and without these corrections shown in Fig. 1.

Models of quark compositeness are probed by searching for evidence of new interactions between quarks at a large characteristic energy scale, Λ . At energies below this scale, the details of the new interaction and potential mediating particles can be integrated out to form a four-fermion contact interaction model [5,6] described by an effective field theory:

$$L_{qq} = \frac{2\pi}{\Lambda^2} [\eta_{LL}(\bar{q}_L \gamma^\mu q_L)(\bar{q}_L \gamma^\mu q_L) + \eta_{RR}(\bar{q}_R \gamma^\mu q_R)(\bar{q}_R \gamma^\mu q_R) + 2\eta_{RL}(\bar{q}_R \gamma^\mu q_R)(\bar{q}_L \gamma^\mu q_L)], \quad (1)$$

where the quark fields have L and R chiral projections and the coefficients η_{LL} , η_{RR} , and η_{RL} turn on and off various interactions.

In this Letter, a contact interaction (CI) model with a left-chiral color-singlet coupling ($\eta_{LL} = \pm 1$) is used as a benchmark model, as many other models of new phenomena have similar predictions for the dijet scattering angle χ at large m_{jj} . Interference of the signal model with the SM process $q\bar{q} \rightarrow q\bar{q}$ is also included.

Event samples were simulated with both QCD and contact interactions, taking interference into account and using the same event generator, underlying-event tune, and PDF as for the SM simulations. Events were generated for both constructive and destructive interference with $\Lambda = 7$ TeV and $\Lambda = 10$ TeV. The $\Lambda = 7$ TeV sample is then used for extrapolation to other values of Λ , using the fact that the interference term is proportional to $1/\Lambda^2$ and the pure CI cross section is proportional to $1/\Lambda^4$. This procedure is validated with the $\Lambda = 10$ TeV sample. As with the QCD prediction, a K -factor correction is computed to correct the PYTHIA LO + shower prediction to a NLO calculation. Calculations at NLO are provided by CIJET [35] v1.0.

Uncertainties in the SM and signal predictions include theoretical uncertainties and experimental uncertainties on the measured angular distributions. Theoretical uncertainties in the SM and signal predictions are due to the choice of PDF, renormalization and factorization scales, choice of event generator, as well as statistical uncertainties due to limited simulation sample sizes. The impact of the uncertainty in the PDF is estimated using NLOJET++ with three different PDFs: CT10, MSTW2008 [36], and NNPDF23 [37]. These uncertainties are negligible ($< 1\%$), as the choice of PDF largely impacts the total cross section rather than the angular distributions. The uncertainty due to the choice of renormalization and factorization scales were estimated by varying those independently up and down by a factor two in NLOJET++. The resulting uncertainty varies with m_{jj} and χ , rising to 4% at the smallest χ values at high m_{jj} . The uncertainty due to the choice of generator is

estimated by comparing the predictions from the NLO generator POWHEG [38] v1.0 with those of PYTHIA8 with K factors applied. The largest uncertainty due to choice of generator is at the lowest m_{jj} values, where it approaches 20%, while for the highest m_{jj} values and smallest χ , it ranges from 10% to 14%. The uncertainty due to the choice of the showering model is estimated through comparison of POWHEG samples showered and hadronized with PYTHIA8 v8.175 to HERWIG [39] v6.520.2 samples using JIMMY [40,41] v4.31. The largest value of this uncertainty is less than 1% at the highest m_{jj} values and smallest values of χ . Finally the statistical uncertainties on the K factors due to limited simulation sample size are small and set to 1%.

The experimental uncertainty is dominated by the η dependence of the jet energy scale calibration. This uncertainty varies from approximately 15% at small values of χ for the highest m_{jj} values, to a few percent at lower m_{jj} values and higher χ values. The uncertainty in the beam energy is found to introduce a negligible contribution. The total uncertainty at the lowest χ , highest m_{jj} amounts to 20%, decreasing to a few percent at high χ . The total theoretical and experimental uncertainties are shown in Fig. 1.

The p value for the SM hypothesis is (0.25) 0.30 for the (second) highest m_{jj} bin. In the absence of significant deviations from the SM prediction, upper bounds on CI contributions are calculated using a one-sided profile likelihood ratio and the CL_s technique [42,43], evaluated using the asymptotic approximation [44] on events with $m_{jj} > 3.2$ TeV; the validity of asymptotic approximation was confirmed using toy simulations. These bounds exclude a compositeness scale below 8.1 TeV in a destructive interference scenario and below 12.0 TeV in a constructive interference scenario. The median expected limits are 8.9 (14.1) TeV for the destructive (constructive) interference scenario.

In summary, dijet angular distributions have been measured by the ATLAS experiment in 17.3 fb^{-1} of 8 TeV pp collisions at the LHC. Over a wide angular range and dijet invariant mass spectrum, the data are well described by QCD predictions at NLO. A model of quark compositeness is used as a benchmark for theories of new phenomena that include new forces and mediating particles; such theories predict deviations at small values of χ . A compositeness scale below 8.1 (12.0) TeV in a destructive (constructive) interference scenario is excluded at 95% confidence level, similar to results from the CMS Collaboration [17] and representing a significant enhancement in sensitivity relative to the previous limit (at 7.6 TeV for destructive interference) from the ATLAS Collaboration [11].

We would like to thank S. Dittmaier and A. Huss for providing us with the electroweak correction factors. We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We

acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWFW and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF, DNSRC and Lundbeck Foundation, Denmark; EPLANET, ERC and NSRF, European Union; IN2P3-CNRS, CEA-DSM/IRFU, France; GNSF, Georgia; BMBF, DFG, HGF, MPG and AvH Foundation, Germany; GSRT and NSRF, Greece; RGC, Hong Kong SAR, China; ISF, MINERVA, GIF, I-CORE and Benoziyo Center, Israel; INFN, Italy; MEXT and JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; BRF and RCN, Norway; MNiSW and NCN, Poland; GRICES and FCT, Portugal; MNE/IFA, Romania; MES of Russia and NRC KI, Russian Federation; JINR; MSTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SER, SNSF and Cantons of Bern and Geneva, Switzerland; NSC, Taiwan; TAEK, Turkey; STFC, the Royal Society and Leverhulme Trust, United Kingdom; DOE and NSF, United States of America. The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (USA) and in the Tier-2 facilities worldwide.

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G. Aad,⁸⁵ B. Abbott,¹¹³ J. Abdallah,¹⁵¹ O. Abdinov,¹¹ R. Aben,¹⁰⁷ M. Abolins,⁹⁰ O. S. AbouZeid,¹⁵⁸ H. Abramowicz,¹⁵³ H. Abreu,¹⁵² R. Abreu,³⁰ Y. Abulaiti,^{146a,146b} B. S. Acharya,^{164a,164b} L. Adamczyk,^{38a} D. L. Adams,²⁵ J. Adelman,¹⁰⁸ S. Adomeit,¹⁰⁰ T. Adye,¹³¹ A. A. Affolder,⁷⁴ T. Agatonovic-Jovin,¹³ J. A. Aguilar-Saavedra,^{126a,126f} S. P. Ahlen,²² F. Ahmadov,^{65,c} G. Aielli,^{133a,133b} H. Akerstedt,^{146a,146b} T. P. A. Åkesson,⁸¹ G. Akimoto,¹⁵⁵ A. V. Akimov,⁹⁶ G. L. Alberghi,^{20a,20b} J. Albert,¹⁶⁹ S. Albrand,⁵⁵ M. J. Alconada Verzini,⁷¹ M. Aleksa,³⁰ I. N. Aleksandrov,⁶⁵ C. Alexa,^{26a} G. Alexander,¹⁵³ T. Alexopoulos,¹⁰ M. Alhroob,¹¹³ G. Alimonti,^{91a} L. Alio,⁸⁵ J. Alison,³¹ S. P. Alkire,³⁵ B. M. M. Allbrooke,¹⁸ P. P. Allport,⁷⁴ A. Aloisio,^{104a,104b} A. Alonso,³⁶ F. Alonso,⁷¹ C. Alpigiani,⁷⁶ A. Altheimer,³⁵ B. Alvarez Gonzalez,³⁰ D. Álvarez Piqueras,¹⁶⁷ M. G. Alviggi,^{104a,104b} B. T. Amadio,¹⁵ K. Amako,⁶⁶ Y. Amaral Coutinho,^{24a} C. Amelung,²³ D. Amidei,⁸⁹ S. P. Amor Dos Santos,^{126a,126c} A. Amorim,^{126a,126b} S. Amoroso,⁴⁸ N. Amram,¹⁵³ G. Amundsen,²³ C. Anastopoulos,¹³⁹ L. S. Ancu,⁴⁹ N. Andari,³⁰ T. Andeen,³⁵ C. F. Anders,^{58b} G. Anders,³⁰ J. K. Anders,⁷⁴ K. J. Anderson,³¹ A. Andreazza,^{91a,91b} V. Andrei,^{58a} S. Angelidakis,⁹ I. Angelozzi,¹⁰⁷ P. Anger,⁴⁴ A. Angerami,³⁵ F. Anghinolfi,³⁰ A. V. Anisenkov,^{109,d} N. Anjos,¹² A. Annovi,^{124a,124b} M. Antonelli,⁴⁷ A. Antonov,⁹⁸ J. Antos,^{144b} F. Anulli,^{132a} M. Aoki,⁶⁶ L. Aperio Bella,¹⁸ G. Arabidze,⁹⁰ Y. Arai,⁶⁶ J. P. Araque,^{126a} A. T. H. Arce,⁴⁵ F. A. Arduh,⁷¹ J-F. Arguin,⁹⁵ S. Argyropoulos,⁴² M. Arik,^{19a} A. J. Armbruster,³⁰ O. Arnaez,³⁰ V. Arnal,⁸² H. Arnold,⁴⁸ M. Arratia,²⁸ O. Arslan,²¹ A. Artamonov,⁹⁷ G. Artoni,²³ S. Asai,¹⁵⁵ N. Asbah,⁴² A. Ashkenazi,¹⁵³ B. Åsman,^{146a,146b} L. Asquith,¹⁴⁹ K. Assamagan,²⁵ R. Astalos,^{144a} M. Atkinson,¹⁶⁵ N. B. Atlay,¹⁴¹ B. Auerbach,⁶ K. Augsten,¹²⁸ M. Aurousseau,^{145b} G. Avolio,³⁰ B. Axen,¹⁵ M. K. Ayoub,¹¹⁷ G. Azuelos,^{95,e} M. A. Baak,³⁰ A. E. Baas,^{58a} C. Bacci,^{134a,134b} H. Bachacou,¹³⁶ K. Bachas,¹⁵⁴ M. Backes,³⁰ M. Backhaus,³⁰ E. Badescu,^{26a} P. Bagiacchi,^{132a,132b} P. Bagnaia,^{132a,132b} Y. Bai,^{33a} T. Bain,³⁵ J. T. Baines,¹³¹ O. K. Baker,¹⁷⁶ P. Balek,¹²⁹ T. Balestri,¹⁴⁸ F. Balli,⁸⁴ E. Banas,³⁹ Sw. Banerjee,¹⁷³ A. A. E. Bannoura,¹⁷⁵ H. S. Bansil,¹⁸ L. Barak,³⁰ S. P. Baranov,⁹⁶ E. L. Barberio,⁸⁸ D. Barberis,^{50a,50b} M. Barbero,⁸⁵ T. Barillari,¹⁰¹ M. Barisonzi,^{164a,164b} T. Barklow,¹⁴³ N. Barlow,²⁸ S. L. Barnes,⁸⁴ B. M. Barnett,¹³¹ R. M. Barnett,¹⁵ Z. Barnovska,⁵ A. Baroncelli,^{134a} G. Barone,⁴⁹ A. J. Barr,¹²⁰ F. Barreiro,⁸² J. Barreiro Guimarães da Costa,⁵⁷ R. Bartoldus,¹⁴³ A. E. Barton,⁷² P. Bartos,^{144a} A. Bassalat,¹¹⁷ A. Basye,¹⁶⁵ R. L. Bates,⁵³ S. J. Batista,¹⁵⁸ J. R. Batley,²⁸ M. Battaglia,¹³⁷ M. Bauce,^{132a,132b} F. Bauer,¹³⁶ H. S. Bawa,^{143,f} J. B. Beacham,¹¹¹ M. D. Beattie,⁷² T. Beau,⁸⁰ P. H. Beauchemin,¹⁶¹ R. Beccherle,^{124a,124b} P. Bechtel,²¹ H. P. Beck,^{17,g} K. Becker,¹²⁰ M. Becker,⁸³ S. Becker,¹⁰⁰ M. Beckingham,¹⁷⁰ C. Becot,¹¹⁷ A. J. Beddall,^{19c} A. Beddall,^{19c} V. A. Bednyakov,⁶⁵ C. P. Bee,¹⁴⁸ L. J. Beamster,¹⁰⁷ T. A. Beermann,¹⁷⁵ M. Begel,²⁵ J. K. Behr,¹²⁰ C. Belanger-Champagne,⁸⁷ P. J. Bell,⁴⁹ W. H. Bell,⁴⁹ G. Bella,¹⁵³ L. Bellagamba,^{20a} A. Bellerive,²⁹ M. Bellomo,⁸⁶ K. Belotskiy,⁹⁸ O. Beltramello,³⁰ O. Benary,¹⁵³ D. Bencheekroun,^{135a} M. Bender,¹⁰⁰ K. Bendtz,^{146a,146b} N. Benekos,¹⁰ Y. Benhammou,¹⁵³ E. Benhar Noccioli,⁴⁹ J. A. Benitez Garcia,^{159b} D. P. Benjamin,⁴⁵ J. R. Bensinger,²³ S. Bentvelsen,¹⁰⁷ L. Beresford,¹²⁰ M. Beretta,⁴⁷ D. Berge,¹⁰⁷ E. Bergeas Kuutmann,¹⁶⁶ N. Berger,⁵ F. Berghaus,¹⁶⁹ J. Beringer,¹⁵ C. Bernard,²² N. R. Bernard,⁸⁶ C. Bernius,¹¹⁰ F. U. Bernlochner,²¹ T. Berry,⁷⁷ P. Berta,¹²⁹ C. Bertella,⁸³ G. Bertoli,^{146a,146b} F. Bertolucci,^{124a,124b} C. Bertsche,¹¹³ D. Bertsche,¹¹³ M. I. Besana,^{91a} G. J. Besjes,¹⁰⁶ O. Bessidskaia Bylund,^{146a,146b} M. Bessner,⁴² N. Besson,¹³⁶ C. Betancourt,⁴⁸ S. Bethke,¹⁰¹ A. J. Bevan,⁷⁶ W. Bhimji,⁴⁶ R. M. Bianchi,¹²⁵ L. Bianchini,²³ M. Bianco,³⁰ O. Biebel,¹⁰⁰ S. P. Bieniek,⁷⁸ M. Biglietti,^{134a} J. Bilbao De Mendizabal,⁴⁹ H. Bilokon,⁴⁷ M. Bindi,⁵⁴ S. Binet,¹¹⁷ A. Bingul,^{19c} C. Bini,^{132a,132b} C. W. Black,¹⁵⁰ J. E. Black,¹⁴³ K. M. Black,²² D. Blackburn,¹³⁸ R. E. Blair,⁶ J.-B. Blanchard,¹³⁶ J. E. Blanco,⁷⁷ T. Blazek,^{144a} I. Bloch,⁴² C. Blocker,²³ W. Blum,^{83,a} U. Blumenschein,⁵⁴ G. J. Bobbink,¹⁰⁷ V. S. Bobrovnikov,^{109,d} S. S. Bocchetta,⁸¹ A. Bocci,⁴⁵ C. Bock,¹⁰⁰ M. Boehler,⁴⁸ J. A. Bogaerts,³⁰ A. G. Bogdanichikov,¹⁰⁹ C. Boehm,^{146a} V. Boisvert,⁷⁷ T. Bold,^{38a} V. Boldea,^{26a} A. S. Boldyrev,⁹⁹ M. Bomben,⁸⁰ M. Bona,⁷⁶ M. Boonekamp,¹³⁶ A. Borisov,¹³⁰ G. Borissov,⁷² S. Borroni,⁴² J. Bortfeldt,¹⁰⁰ V. Bortolotto,^{60a,60b,60c} K. Bos,¹⁰⁷ D. Boscherini,^{20a} M. Bosman,¹² J. Boudreau,¹²⁵ J. Bouffard,² E. V. Bouhova-Thacker,⁷² D. Boumediene,³⁴ C. Bourdarios,¹¹⁷ N. Bousson,¹¹⁴ A. Boveia,³⁰ J. Boyd,³⁰ I. R. Boyko,⁶⁵ I. Bozic,¹³ J. Bracinik,¹⁸ A. Brandt,⁸ G. Brandt,⁵⁴ O. Brandt,^{58a} U. Bratzler,¹⁵⁶ B. Brau,⁸⁶ J. E. Brau,¹¹⁶ H. M. Braun,^{175,a} S. F. Brazzale,^{164a,164c} K. Brendlinger,¹²² A. J. Brennan,⁸⁸ L. Brenner,¹⁰⁷ R. Brenner,¹⁶⁶ S. Bressler,¹⁷² K. Bristow,^{145c} T. M. Bristow,⁴⁶ D. Britton,⁵³ D. Britzger,⁴² F. M. Brochu,²⁸ I. Brock,²¹ R. Brock,⁹⁰ J. Bronner,¹⁰¹ G. Brooijmans,³⁵ T. Brooks,⁷⁷ W. K. Brooks,^{32b} J. Brosamer,¹⁵ E. Brost,¹¹⁶ J. Brown,⁵⁵ P. A. Bruckman de Renstrom,³⁹ D. Bruncko,^{144b} R. Bruneliere,⁴⁸ A. Bruni,^{20a} G. Bruni,^{20a} M. Bruschi,^{20a} L. Bryngemark,⁸¹ T. Buanes,¹⁴ Q. Buat,¹⁴² P. Buchholz,¹⁴¹ A. G. Buckley,⁵³ S. I. Buda,^{26a} I. A. Budagov,⁶⁵ F. Buehrer,⁴⁸ L. Bugge,¹¹⁹ M. K. Bugge,¹¹⁹ O. Bulekov,⁹⁸ D. Bullock,⁸ H. Burckhart,³⁰ S. Burdin,⁷⁴ B. Burghgrave,¹⁰⁸ S. Burke,¹³¹ I. Burmeister,⁴³ E. Busato,³⁴ D. Büscher,⁴⁸ V. Büscher,⁸³ P. Bussey,⁵³ C. P. Buszello,¹⁶⁶ J. M. Butler,²² A. I. Butt,³ C. M. Buttar,⁵³ J. M. Butterworth,⁷⁸ P. Butti,¹⁰⁷ W. Buttinger,²⁵ A. Buzatu,⁵³ R. Buzykaev,^{109,d} S. Cabrera Urbán,¹⁶⁷ D. Caforio,¹²⁸

V. M. Cairo,^{37a,37b} O. Cakir,^{4a} P. Calafiura,¹⁵ A. Calandri,¹³⁶ G. Calderini,⁸⁰ P. Calfayan,¹⁰⁰ L. P. Caloba,^{24a} D. Calvet,³⁴ S. Calvet,³⁴ R. Camacho Toro,³¹ S. Camarda,⁴² P. Camarri,^{133a,133b} D. Cameron,¹¹⁹ L. M. Caminada,¹⁵ R. Caminal Armadans,¹² S. Campana,³⁰ M. Campanelli,⁷⁸ A. Campoverde,¹⁴⁸ V. Canale,^{104a,104b} A. Canepa,^{159a} M. Cano Bret,⁷⁶ J. Cantero,⁸² R. Cantrill,^{126a} T. Cao,⁴⁰ M. D. M. Capeans Garrido,³⁰ I. Caprini,^{26a} M. Caprini,^{26a} M. Capua,^{37a,37b} R. Caputo,⁸³ R. Cardarelli,^{133a} T. Carli,³⁰ G. Carlino,^{104a} L. Carminati,^{91a,91b} S. Caron,¹⁰⁶ E. Carquin,^{32a} G. D. Carrillo-Montoya,⁸ J. R. Carter,²⁸ J. Carvalho,^{126a,126c} D. Casadei,⁷⁸ M. P. Casado,¹² M. Casolino,¹² E. Castaneda-Miranda,^{145b} A. Castelli,¹⁰⁷ V. Castillo Gimenez,¹⁶⁷ N. F. Castro,^{126a,h} P. Catastini,⁵⁷ A. Catinaccio,³⁰ J. R. Catmore,¹¹⁹ A. Cattai,³⁰ J. Caudron,⁸³ V. Cavaliere,¹⁶⁵ D. Cavalli,^{91a} M. Cavalli-Sforza,¹² V. Cavasinni,^{124a,124b} F. Ceradini,^{134a,134b} B. C. Cerio,⁴⁵ K. Cerny,¹²⁹ A. S. Cerqueira,^{24b} A. Cerri,¹⁴⁹ L. Cerrito,⁷⁶ F. Cerutti,¹⁵ M. Cerv,³⁰ A. Cervelli,¹⁷ S. A. Cetin,^{19b} A. Chafaq,^{135a} D. Chakraborty,¹⁰⁸ I. Chalupkova,¹²⁹ P. Chang,¹⁶⁵ B. Chapleau,⁸⁷ J. D. Chapman,²⁸ D. G. Charlton,¹⁸ C. C. Chau,¹⁵⁸ C. A. Chavez Barajas,¹⁴⁹ S. Cheatham,¹⁵² A. Chegwidden,⁹⁰ S. Chekanov,⁶ S. V. Chekulaev,^{159a} G. A. Chelkov,^{65,i} M. A. Chelstowska,⁸⁹ C. Chen,⁶⁴ H. Chen,²⁵ K. Chen,¹⁴⁸ L. Chen,^{33d,j} S. Chen,^{33c} X. Chen,^{33f} Y. Chen,⁶⁷ H. C. Cheng,⁸⁹ Y. Cheng,³¹ A. Cheplakov,⁶⁵ E. Cheremushkina,¹³⁰ R. Cherkaoui El Moursli,^{135e} V. Chernyatin,^{25,a} E. Cheu,⁷ L. Chevalier,¹³⁶ V. Chiarella,⁴⁷ J. T. Childers,⁶ G. Chiodini,^{73a} A. S. Chisholm,¹⁸ R. T. Chislett,⁷⁸ A. Chitan,^{26a} M. V. Chizhov,⁶⁵ K. Choi,⁶¹ S. Chouridou,⁹ B. K. B. Chow,¹⁰⁰ V. Christodoulou,⁷⁸ D. Chromek-Burckhart,³⁰ M. L. Chu,¹⁵¹ J. Chudoba,¹²⁷ A. J. Chuinard,⁸⁷ J. J. Chwastowski,³⁹ L. Chytka,¹¹⁵ G. Ciapetti,^{132a,132b} A. K. Ciftci,^{4a} D. Cinca,⁵³ V. Cindro,⁷⁵ I. A. Cioara,²¹ A. Ciocio,¹⁵ Z. H. Citron,¹⁷² M. Ciubancan,^{26a} A. Clark,⁴⁹ B. L. Clark,⁵⁷ P. J. Clark,⁴⁶ R. N. Clarke,¹⁵ W. Cleland,¹²⁵ C. Clement,^{146a,146b} Y. Coadou,⁸⁵ M. Cobal,^{164a,164c} A. Cocco,¹³⁸ J. Cochran,⁶⁴ L. Coffey,²³ J. G. Cogan,¹⁴³ B. Cole,³⁵ S. Cole,¹⁰⁸ A. P. Colijn,¹⁰⁷ J. Collot,⁵⁵ T. Colombo,^{58c} G. Compostella,¹⁰¹ P. Conde Muiño,^{126a,126b} E. Coniavitis,⁴⁸ S. H. Connell,^{145b} I. A. Connelly,⁷⁷ S. M. Consonni,^{91a,91b} V. Consorti,⁴⁸ S. Constantinescu,^{26a} C. Conta,^{121a,121b} G. Conti,³⁰ F. Conventi,^{104a,k} M. Cooke,¹⁵ B. D. Cooper,⁷⁸ A. M. Cooper-Sarkar,¹²⁰ T. Cornelissen,¹⁷⁵ M. Corradi,^{20a} F. Corriveau,^{87,l} A. Corso-Radu,¹⁶³ A. Cortes-Gonzalez,¹² G. Cortiana,¹⁰¹ G. Costa,^{91a} M. J. Costa,¹⁶⁷ D. Costanzo,¹³⁹ D. Côté,⁸ G. Cottin,²⁸ G. Cowan,⁷⁷ B. E. Cox,⁸⁴ K. Cranmer,¹¹⁰ G. Cree,²⁹ S. Crépe-Renaudin,⁵⁵ F. Crescioli,⁸⁰ W. A. Cribbs,^{146a,146b} M. Crispin Ortuzar,¹²⁰ M. Cristinziani,²¹ V. Croft,¹⁰⁶ G. Crosetti,^{37a,37b} T. Cuhadar Donszelmann,¹³⁹ J. Cummings,¹⁷⁶ M. Curatolo,⁴⁷ C. Cuthbert,¹⁵⁰ H. Czirr,¹⁴¹ P. Czodrowski,³ S. D'Auria,⁵³ M. D'Onofrio,⁷⁴ M. J. Da Cunha Sargedas De Sousa,^{126a,126b} C. Da Via,⁸⁴ W. Dabrowski,^{38a} A. Dafinca,¹²⁰ T. Dai,⁸⁹ O. Dale,¹⁴ F. Dallaire,⁹⁵ C. Dallapiccola,⁸⁶ M. Dam,³⁶ J. R. Dandoy,³¹ N. P. Dang,⁴⁸ A. C. Daniells,¹⁸ M. Danninger,¹⁶⁸ M. Dano Hoffmann,¹³⁶ V. Dao,⁴⁸ G. Darbo,^{50a} S. Darmora,⁸ J. Dassoulas,³ A. Dattagupta,⁶¹ W. Davey,²¹ C. David,¹⁶⁹ T. Davidek,¹²⁹ E. Davies,^{120,m} M. Davies,¹⁵³ P. Davison,⁷⁸ Y. Davygora,^{58a} E. Dawe,⁸⁸ I. Dawson,¹³⁹ R. K. Daya-Ishmukhametova,⁸⁶ K. De,⁸ R. de Asmundis,^{104a} S. De Castro,^{20a,20b} S. De Cecco,⁸⁰ N. De Groot,¹⁰⁶ P. de Jong,¹⁰⁷ H. De la Torre,⁸² F. De Lorenzi,⁶⁴ L. De Nooij,¹⁰⁷ D. De Pedis,^{132a} A. De Salvo,^{132a} U. De Sanctis,¹⁴⁹ A. De Santo,¹⁴⁹ J. B. De Vivie De Regie,¹¹⁷ W. J. Dearnaley,⁷² R. Debbe,²⁵ C. Debenedetti,¹³⁷ D. V. Dedovich,⁶⁵ I. Deigaard,¹⁰⁷ J. Del Peso,⁸² T. Del Prete,^{124a,124b} D. Delgove,¹¹⁷ F. Deliot,¹³⁶ C. M. Delitzsch,⁴⁹ M. Deliyergiyev,⁷⁵ A. Dell'Acqua,³⁰ L. Dell'Asta,²² M. Dell'Orso,^{124a,124b} M. Della Pietra,^{104a,k} D. della Volpe,⁴⁹ M. Delmastro,⁵ P. A. Delsart,⁵⁵ C. Deluca,¹⁰⁷ D. A. DeMarco,¹⁵⁸ S. Demers,¹⁷⁶ M. Demichev,⁶⁵ A. Demilly,⁸⁰ S. P. Denisov,¹³⁰ D. Derendarz,³⁹ J. E. Derkaoui,^{135d} F. Derue,⁸⁰ P. Dervan,⁷⁴ K. Desch,²¹ C. Deterre,⁴² P. O. Deviveiros,³⁰ A. Dewhurst,¹³¹ S. Dhaliwal,¹⁰⁷ A. Di Ciaccio,^{133a,133b} L. Di Ciaccio,⁵ A. Di Domenico,^{132a,132b} C. Di Donato,^{104a,104b} A. Di Girolamo,³⁰ B. Di Girolamo,³⁰ A. Di Mattia,¹⁵² B. Di Micco,^{134a,134b} R. Di Nardo,⁴⁷ A. Di Simone,⁴⁸ R. Di Sipio,¹⁵⁸ D. Di Valentino,²⁹ C. Diaconu,⁸⁵ M. Diamond,¹⁵⁸ F. A. Dias,⁴⁶ M. A. Diaz,^{32a} E. B. Diehl,⁸⁹ J. Dietrich,¹⁶ S. Diglio,⁸⁵ A. Dimitrievska,¹³ J. Dingfelder,²¹ F. Dittus,³⁰ F. Djama,⁸⁵ T. Djobava,^{51b} J. I. Djuvsland,^{58a} M. A. B. do Vale,^{24c} D. Dobos,³⁰ M. Dobre,^{26a} C. Doglioni,⁴⁹ T. Dohmae,¹⁵⁵ J. Dolejsi,¹²⁹ Z. Dolezal,¹²⁹ B. A. Dolgoshein,^{98,a} M. Donadelli,^{24d} S. Donati,^{124a,124b} P. Dondero,^{121a,121b} J. Donini,³⁴ J. Dopke,¹³¹ A. Doria,^{104a} M. T. Dova,⁷¹ A. T. Doyle,⁵³ E. Drechsler,⁵⁴ M. Dris,¹⁰ E. Dubreuil,³⁴ E. Duchovni,¹⁷² G. Duckeck,¹⁰⁰ O. A. Ducu,^{26a,85} D. Duda,¹⁷⁵ A. Dudarev,³⁰ L. Duflot,¹¹⁷ L. Duguid,⁷⁷ M. Dührssen,³⁰ M. Dunford,^{58a} H. Duran Yildiz,^{4a} M. Düren,⁵² A. Durglishvili,^{51b} D. Duschinger,⁴⁴ M. Dyndal,^{38a} C. Eckardt,⁴² K. M. Ecker,¹⁰¹ R. C. Edgar,⁸⁹ W. Edson,² N. C. Edwards,⁴⁶ W. Ehrenfeld,²¹ T. Eifert,³⁰ G. Eigen,¹⁴ K. Einsweiler,¹⁵ T. Ekelof,¹⁶⁶ M. El Kacimi,^{135c} M. Ellert,¹⁶⁶ S. Elles,⁵ F. Ellinghaus,⁸³ A. A. Elliot,¹⁶⁹ N. Ellis,³⁰ J. Elmsheuser,¹⁰⁰ M. Elsing,³⁰ D. Emeliyanov,¹³¹ Y. Enari,¹⁵⁵ O. C. Endner,⁸³ M. Endo,¹¹⁸ R. Engelmann,¹⁴⁸ J. Erdmann,⁴³ A. Ereditato,¹⁷ G. Ernis,¹⁷⁵ J. Ernst,² M. Ernst,²⁵ S. Errede,¹⁶⁵ E. Ertel,⁸³ M. Escalier,¹¹⁷ H. Esch,⁴³ C. Escobar,¹²⁵ B. Esposito,⁴⁷ A. I. Etienvre,¹³⁶ E. Etzion,¹⁵³ H. Evans,⁶¹ A. Ezhilov,¹²³ L. Fabbri,^{20a,20b} G. Facini,³¹ R. M. Fakhruddinov,¹³⁰

S. Falciano,^{132a} R. J. Falla,⁷⁸ J. Faltova,¹²⁹ Y. Fang,^{33a} M. Fantì,^{91a,91b} A. Farbin,⁸ A. Farilla,^{134a} T. Farooque,¹² S. Farrell,¹⁵ S. M. Farrington,¹⁷⁰ P. Farthouat,³⁰ F. Fassi,^{135e} P. Fassnacht,³⁰ D. Fassouliotis,⁹ M. Faucci Giannelli,⁷⁷ A. Favareto,^{50a,50b} L. Fayard,¹¹⁷ P. Federic,^{144a} O. L. Fedin,^{123,n} W. Fedorko,¹⁶⁸ S. Feigl,³⁰ L. Feligioni,⁸⁵ C. Feng,^{33d} E. J. Feng,⁶ H. Feng,⁸⁹ A. B. Fenyuk,¹³⁰ P. Fernandez Martinez,¹⁶⁷ S. Fernandez Perez,³⁰ S. Ferrag,⁵³ J. Ferrando,⁵³ A. Ferrari,¹⁶⁶ P. Ferrari,¹⁰⁷ R. Ferrari,^{121a} D. E. Ferreira de Lima,⁵³ A. Ferrer,¹⁶⁷ D. Ferrere,⁴⁹ C. Ferretti,⁸⁹ A. Ferretto Parodi,^{50a,50b} M. Fiascaris,³¹ F. Fiedler,⁸³ A. Filipčič,⁷⁵ M. Filipuzzi,⁴² F. Filthaut,¹⁰⁶ M. Fincke-Keeler,¹⁶⁹ K. D. Finelli,¹⁵⁰ M. C. N. Fiolhais,^{126a,126c} L. Fiorini,¹⁶⁷ A. Firan,⁴⁰ A. Fischer,² C. Fischer,¹² J. Fischer,¹⁷⁵ W. C. Fisher,⁹⁰ E. A. Fitzgerald,²³ M. Flechl,⁴⁸ I. Fleck,¹⁴¹ P. Fleischmann,⁸⁹ S. Fleischmann,¹⁷⁵ G. T. Fletcher,¹³⁹ G. Fletcher,⁷⁶ T. Flick,¹⁷⁵ A. Floderus,⁸¹ L. R. Flores Castillo,^{60a} M. J. Flowerdew,¹⁰¹ A. Formica,¹³⁶ A. Forti,⁸⁴ D. Fournier,¹¹⁷ H. Fox,⁷² S. Fracchia,¹² P. Francavilla,⁸⁰ M. Franchini,^{20a,20b} D. Francis,³⁰ L. Franconi,¹¹⁹ M. Franklin,⁵⁷ M. Fraternali,^{121a,121b} D. Freeborn,⁷⁸ S. T. French,²⁸ F. Friedrich,⁴⁴ D. Froidevaux,³⁰ J. A. Frost,¹²⁰ C. Fukunaga,¹⁵⁶ E. Fullana Torregrosa,⁸³ B. G. Fulson,¹⁴³ J. Fuster,¹⁶⁷ C. Gabaldon,⁵⁵ O. Gabizon,¹⁷⁵ A. Gabrielli,^{20a,20b} A. Gabrielli,^{132a,132b} S. Gadatsch,¹⁰⁷ S. Gadomski,⁴⁹ G. Gagliardi,^{50a,50b} P. Gagnon,⁶¹ C. Galea,¹⁰⁶ B. Galhardo,^{126a,126c} E. J. Gallas,¹²⁰ B. J. Gallop,¹³¹ P. Gallus,¹²⁸ G. Galster,³⁶ K. K. Gan,¹¹¹ J. Gao,^{33b,85} Y. Gao,⁴⁶ Y. S. Gao,^{143,f} F. M. Garay Walls,⁴⁶ F. Garberson,¹⁷⁶ C. García,¹⁶⁷ J. E. García Navarro,¹⁶⁷ M. Garcia-Sciveres,¹⁵ R. W. Gardner,³¹ N. Garelli,¹⁴³ V. Garonne,¹¹⁹ C. Gatti,⁴⁷ A. Gaudiello,^{50a,50b} G. Gaudio,^{121a} B. Gaur,¹⁴¹ L. Gauthier,⁹⁵ P. Gauzzi,^{132a,132b} I. L. Gavrilenko,⁹⁶ C. Gay,¹⁶⁸ G. Gaycken,²¹ E. N. Gazis,¹⁰ P. Ge,^{33d} Z. Gecse,¹⁶⁸ C. N. P. Gee,¹³¹ D. A. A. Geerts,¹⁰⁷ Ch. Geich-Gimbel,²¹ M. P. Geisler,^{58a} C. Gemme,^{50a} M. H. Genest,⁵⁵ S. Gentile,^{132a,132b} M. George,⁵⁴ S. George,⁷⁷ D. Gerbaudo,¹⁶³ A. Gershon,¹⁵³ H. Ghazlane,^{135b} B. Giacobbe,^{20a} S. Giagu,^{132a,132b} V. Giangiobbe,¹² P. Giannetti,^{124a,124b} B. Gibbard,²⁵ S. M. Gibson,⁷⁷ M. Gilchriese,¹⁵ T. P. S. Gillam,²⁸ D. Gillberg,³⁰ G. Gilles,³⁴ D. M. Gingrich,^{3,e} N. Giokaris,⁹ M. P. Giordani,^{164a,164c} F. M. Giorgi,^{20a} F. M. Giorgi,¹⁶ P. F. Giraud,¹³⁶ P. Giromini,⁴⁷ D. Giugni,^{91a} C. Giuliani,⁴⁸ M. Giulini,^{58b} B. K. Gjelsten,¹¹⁹ S. Gkaitatzis,¹⁵⁴ I. Gkialas,¹⁵⁴ E. L. Gkoukousis,¹¹⁷ L. K. Gladilin,⁹⁹ C. Glasman,⁸² J. Glatzer,³⁰ P. C. F. Glaysher,⁴⁶ A. Glazov,⁴² M. Goblirsch-Kolb,¹⁰¹ J. R. Goddard,⁷⁶ J. Godlewski,³⁹ S. Goldfarb,⁸⁹ T. Golling,⁴⁹ D. Golubkov,¹³⁰ A. Gomes,^{126a,126b,126d} R. Gonçalves,^{126a} J. Goncalves Pinto Firmino Da Costa,¹³⁶ L. Gonella,²¹ S. González de la Hoz,¹⁶⁷ G. Gonzalez Parra,¹² S. Gonzalez-Sevilla,⁴⁹ L. Goossens,³⁰ P. A. Gorbounov,⁹⁷ H. A. Gordon,²⁵ I. Gorelov,¹⁰⁵ B. Gorini,³⁰ E. Gorini,^{73a,73b} A. Gorišek,⁷⁵ E. Gornicki,³⁹ A. T. Goshaw,⁴⁵ C. Gössling,⁴³ M. I. Gostkin,⁶⁵ D. Goujdami,^{135c} A. G. Goussiou,¹³⁸ N. Govender,^{145b} H. M. X. Grabas,¹³⁷ L. Graber,⁵⁴ I. Grabowska-Bold,^{38a} P. Grafström,^{20a,20b} K.-J. Grahm,⁴² J. Gramling,⁴⁹ E. Gramstad,¹¹⁹ S. Grancagnolo,¹⁶ V. Grassi,¹⁴⁸ V. Gratchev,¹²³ H. M. Gray,³⁰ E. Graziani,^{134a} Z. D. Greenwood,^{79,o} K. Gregersen,⁷⁸ I. M. Gregor,⁴² P. Grenier,¹⁴³ J. Griffiths,⁸ A. A. Grillo,¹³⁷ K. Grimm,⁷² S. Grinstein,^{12,p} Ph. Gris,³⁴ J.-F. Grivaz,¹¹⁷ J. P. Grohs,⁴⁴ A. Grohsjean,⁴² E. Gross,¹⁷² J. Grosse-Knetter,⁵⁴ G. C. Grossi,⁷⁹ Z. J. Grout,¹⁴⁹ L. Guan,^{33b} J. Guenther,¹²⁸ F. Guescini,⁴⁹ D. Guest,¹⁷⁶ O. Gueta,¹⁵³ E. Guido,^{50a,50b} T. Guillemin,¹¹⁷ S. Guindon,² U. Gul,⁵³ C. Gumpert,⁴⁴ J. Guo,^{33e} S. Gupta,¹²⁰ P. Gutierrez,¹¹³ N. G. Gutierrez Ortiz,⁵³ C. Gutsche,⁴⁴ C. Guyot,¹³⁶ C. Gwenlan,¹²⁰ C. B. Gwilliam,⁷⁴ A. Haas,¹¹⁰ C. Haber,¹⁵ H. K. Hadavand,⁸ N. Haddad,^{135e} P. Haefner,²¹ S. Hageböck,²¹ Z. Hajduk,³⁹ H. Hakobyan,¹⁷⁷ M. Haleem,⁴² J. Haley,¹¹⁴ D. Hall,¹²⁰ G. Halladjian,⁹⁰ G. D. Hallewell,⁸⁵ K. Hamacher,¹⁷⁵ P. Hamal,¹¹⁵ K. Hamano,¹⁶⁹ M. Hamer,⁵⁴ A. Hamilton,^{145a} G. N. Hamity,^{145c} P. G. Hamnett,⁴² L. Han,^{33b} K. Hanagaki,¹¹⁸ K. Hanawa,¹⁵⁵ M. Hance,¹⁵ P. Hanke,^{58a} R. Hanna,¹³⁶ J. B. Hansen,³⁶ J. D. Hansen,³⁶ M. C. Hansen,²¹ P. H. Hansen,³⁶ K. Hara,¹⁶⁰ A. S. Hard,¹⁷³ T. Harenberg,¹⁷⁵ F. Hariri,¹¹⁷ S. Harkusha,⁹² R. D. Harrington,⁴⁶ P. F. Harrison,¹⁷⁰ F. Hartjes,¹⁰⁷ M. Hasegawa,⁶⁷ S. Hasegawa,¹⁰³ Y. Hasegawa,¹⁴⁰ A. Hasib,¹¹³ S. Hassani,¹³⁶ S. Haug,¹⁷ R. Hauser,⁹⁰ L. Hauswald,⁴⁴ M. Havranek,¹²⁷ C. M. Hawkes,¹⁸ R. J. Hawkins,³⁰ A. D. Hawkins,⁸¹ T. Hayashi,¹⁶⁰ D. Hayden,⁹⁰ C. P. Hays,¹²⁰ J. M. Hays,⁷⁶ H. S. Hayward,⁷⁴ S. J. Haywood,¹³¹ S. J. Head,¹⁸ T. Heck,⁸³ V. Hedberg,⁸¹ L. Heelan,⁸ S. Heim,¹²² T. Heim,¹⁷⁵ B. Heinemann,¹⁵ L. Heinrich,¹¹⁰ J. Hejbal,¹²⁷ L. Helary,²² S. Hellman,^{146a,146b} D. Hellmich,²¹ C. Helsens,³⁰ J. Henderson,¹²⁰ R. C. W. Henderson,⁷² Y. Heng,¹⁷³ C. Hengler,⁴² A. Henrichs,¹⁷⁶ A. M. Henriques Correia,³⁰ S. Henrot-Versille,¹¹⁷ G. H. Herbert,¹⁶ Y. Hernández Jiménez,¹⁶⁷ R. Herrberg-Schubert,¹⁶ G. Herten,⁴⁸ R. Hertenberger,¹⁰⁰ L. Hervas,³⁰ G. G. Hesketh,⁷⁸ N. P. Hesse,¹⁰⁷ J. W. Hetherly,⁴⁰ R. Hickling,⁷⁶ E. Higón-Rodríguez,¹⁶⁷ E. Hill,¹⁶⁹ J. C. Hill,²⁸ K. H. Hiller,⁴² S. J. Hillier,¹⁸ I. Hinchliffe,¹⁵ E. Hines,¹²² R. R. Hinman,¹⁵ M. Hirose,¹⁵⁷ D. Hirschbuehl,¹⁷⁵ J. Hobbs,¹⁴⁸ N. Hod,¹⁰⁷ M. C. Hodgkinson,¹³⁹ P. Hodgson,¹³⁹ A. Hoecker,³⁰ M. R. Hoferkamp,¹⁰⁵ F. Hoenig,¹⁰⁰ M. Hohlfeld,⁸³ D. Hohn,²¹ T. R. Holmes,¹⁵ T. M. Hong,¹²² L. Hooft van Huysduynen,¹¹⁰ W. H. Hopkins,¹¹⁶ Y. Horii,¹⁰³ A. J. Horton,¹⁴² J.-Y. Hostachy,⁵⁵ S. Hou,¹⁵¹ A. Hoummada,^{135a} J. Howard,¹²⁰ J. Howarth,⁴² M. Hrabovsky,¹¹⁵ I. Hristova,¹⁶ J. Hrivnac,¹¹⁷ T. Hryn'ova,⁵ A. Hrynevich,⁹³ C. Hsu,^{145c} P. J. Hsu,^{151,q} S.-C. Hsu,¹³⁸ D. Hu,³⁵ Q. Hu,^{33b} X. Hu,⁸⁹ Y. Huang,⁴² Z. Hubacek,³⁰

F. Hubaut,⁸⁵ F. Huegging,²¹ T. B. Huffman,¹²⁰ E. W. Hughes,³⁵ G. Hughes,⁷² M. Huhtinen,³⁰ T. A. Hülsing,⁸³
 N. Huseynov,^{65,c} J. Huston,⁹⁰ J. Huth,⁵⁷ G. Iacobucci,⁴⁹ G. Iakovidis,²⁵ I. Ibragimov,¹⁴¹ L. Iconomidou-Fayard,¹¹⁷
 E. Ideal,¹⁷⁶ Z. Idrissi,^{135e} P. Iengo,³⁰ O. Igonkina,¹⁰⁷ T. Iizawa,¹⁷¹ Y. Ikegami,⁶⁶ K. Ikematsu,¹⁴¹ M. Ikeno,⁶⁶ Y. Ilchenko,^{31,r}
 D. Iliadis,¹⁵⁴ N. Ilic,¹⁴³ Y. Inamaru,⁶⁷ T. Ince,¹⁰¹ P. Ioannou,⁹ M. Iodice,^{134a} K. Iordanidou,³⁵ V. Ippolito,⁵⁷ A. Irles Quiles,¹⁶⁷
 C. Isaksson,¹⁶⁶ M. Ishino,⁶⁸ M. Ishitsuka,¹⁵⁷ R. Ishmukhametov,¹¹¹ C. Issever,¹²⁰ S. Istin,^{19a} J. M. Iturbe Ponce,⁸⁴
 R. Iuppa,^{133a,133b} J. Ivarsson,⁸¹ W. Iwanski,³⁹ H. Iwasaki,⁶⁶ J. M. Izen,⁴¹ V. Izzo,^{104a} S. Jabbar,³ B. Jackson,¹²² M. Jackson,⁷⁴
 P. Jackson,¹ M. R. Jaekel,³⁰ V. Jain,² K. Jakobs,⁴⁸ S. Jakobsen,³⁰ T. Jakoubek,¹²⁷ J. Jakubek,¹²⁸ D. O. Jamin,¹⁵¹ D. K. Jana,⁷⁹
 E. Jansen,⁷⁸ R. W. Jansky,⁶² J. Janssen,²¹ M. Janus,¹⁷⁰ G. Jarlskog,⁸¹ N. Javadov,^{65,c} T. Javůrek,⁴⁸ L. Jeanty,¹⁵ J. Jejelava,^{51a,s}
 G.-Y. Jeng,¹⁵⁰ D. Jennens,⁸⁸ P. Jenni,^{48,t} J. Jentzsch,⁴³ C. Jeske,¹⁷⁰ S. Jézéquel,⁵ H. Ji,¹⁷³ J. Jia,¹⁴⁸ Y. Jiang,^{33b} S. Jiggins,⁷⁸
 J. Jimenez Pena,¹⁶⁷ S. Jin,^{33a} A. Jinaru,^{26a} O. Jinnouchi,¹⁵⁷ M. D. Joergensen,³⁶ P. Johansson,¹³⁹ K. A. Johns,⁷
 K. Jon-And,^{146a,146b} G. Jones,¹⁷⁰ R. W. L. Jones,⁷² T. J. Jones,⁷⁴ J. Jongmanns,^{58a} P. M. Jorge,^{126a,126b} K. D. Joshi,⁸⁴
 J. Jovicevic,^{159a} X. Ju,¹⁷³ C. A. Jung,⁴³ P. Jussel,⁶² A. Juste Rozas,^{12,p} M. Kaci,¹⁶⁷ A. Kaczmarska,³⁹ M. Kado,¹¹⁷
 H. Kagan,¹¹¹ M. Kagan,¹⁴³ S. J. Kahn,⁸⁵ E. Kajomovitz,⁴⁵ C. W. Kalderon,¹²⁰ S. Kama,⁴⁰ A. Kamenshchikov,¹³⁰
 N. Kanaya,¹⁵⁵ M. Kaneda,³⁰ S. Kaneti,²⁸ V. A. Kantserov,⁹⁸ J. Kanzaki,⁶⁶ B. Kaplan,¹¹⁰ A. Kapliy,³¹ D. Kar,⁵³
 K. Karakostas,¹⁰ A. Karamaoun,³ N. Karastathis,^{10,107} M. J. Kareem,⁵⁴ M. Karnevskiy,⁸³ S. N. Karpov,⁶⁵ Z. M. Karpova,⁶⁵
 K. Karthik,¹¹⁰ V. Kartvelishvili,⁷² A. N. Karyukhin,¹³⁰ L. Kashif,¹⁷³ R. D. Kass,¹¹¹ A. Kastanas,¹⁴ Y. Kataoka,¹⁵⁵ A. Katre,⁴⁹
 J. Katzy,⁴² K. Kawagoe,⁷⁰ T. Kawamoto,¹⁵⁵ G. Kawamura,⁵⁴ S. Kazama,¹⁵⁵ V. F. Kazanin,^{109,d} M. Y. Kazarinov,⁶⁵
 R. Keeler,¹⁶⁹ R. Kehoe,⁴⁰ J. S. Keller,⁴² J. J. Kempster,⁷⁷ H. Keoshkerian,⁸⁴ O. Kepka,¹²⁷ B. P. Kerševan,⁷⁵ S. Kersten,¹⁷⁵
 R. A. Keyes,⁸⁷ F. Khalil-zada,¹¹ H. Khandanyan,^{146a,146b} A. Khanov,¹¹⁴ A. G. Kharlamov,^{109,d} T. J. Khoo,²⁸ V. Khovanskiy,⁹⁷
 E. Khramov,⁶⁵ J. Khubua,^{51b,u} H. Y. Kim,⁸ H. Kim,^{146a,146b} S. H. Kim,¹⁶⁰ Y. Kim,³¹ N. Kimura,¹⁵⁴ O. M. Kind,¹⁶ B. T. King,⁷⁴
 M. King,¹⁶⁷ R. S. B. King,¹²⁰ S. B. King,¹⁶⁸ J. Kirk,¹³¹ A. E. Kiryunin,¹⁰¹ T. Kishimoto,⁶⁷ D. Kisieleska,^{38a} F. Kiss,⁴⁸
 K. Kiuchi,¹⁶⁰ O. Kivernyk,¹³⁶ E. Kladiva,^{144b} M. H. Klein,³⁵ M. Klein,⁷⁴ U. Klein,⁷⁴ K. Kleinknecht,⁸³ P. Klimek,^{146a,146b}
 A. Klimentov,²⁵ R. Klingenberg,⁴³ J. A. Klinger,⁸⁴ T. Klioutchnikova,³⁰ E.-E. Kluge,^{58a} P. Kluit,¹⁰⁷ S. Kluth,¹⁰¹
 E. Kneringer,⁶² E. B. F. G. Knoops,⁸⁵ A. Knue,⁵³ A. Kobayashi,¹⁵⁵ D. Kobayashi,¹⁵⁷ T. Kobayashi,¹⁵⁵ M. Kobel,⁴⁴
 M. Kocian,¹⁴³ P. Kodys,¹²⁹ T. Koffas,²⁹ E. Koffeman,¹⁰⁷ L. A. Kogan,¹²⁰ S. Kohlmann,¹⁷⁵ Z. Kohout,¹²⁸ T. Kohriki,⁶⁶
 T. Koi,¹⁴³ H. Kolanoski,¹⁶ I. Koletsou,⁵ A. A. Komar,^{96a} Y. Komori,¹⁵⁵ T. Kondo,⁶⁶ N. Kondrashova,⁴² K. Köneke,⁴⁸
 A. C. König,¹⁰⁶ S. König,⁸³ T. Kono,^{66,v} R. Konoplich,^{110,w} N. Konstantinidis,⁷⁸ R. Kopeliansky,¹⁵² S. Koperny,^{38a}
 L. Köpke,⁸³ A. K. Kopp,⁴⁸ K. Korcyl,³⁹ K. Kordas,¹⁵⁴ A. Korn,⁷⁸ A. A. Korol,^{109,d} I. Korolkov,¹² E. V. Korolkova,¹³⁹
 O. Kortner,¹⁰¹ S. Kortner,¹⁰¹ T. Kosek,¹²⁹ V. V. Kostyukhin,²¹ V. M. Kotov,⁶⁵ A. Kotwal,⁴⁵ A. Kourkoumeli-Charalampidi,¹⁵⁴
 C. Kourkoumelis,⁹ V. Kouskoura,²⁵ A. Koutsman,^{159a} R. Kowalewski,¹⁶⁹ T. Z. Kowalski,^{38a} W. Kozanecki,¹³⁶
 A. S. Kozhin,¹³⁰ V. A. Kramarenko,⁹⁹ G. Kramerberger,⁷⁵ D. Krasnopevtsev,⁹⁸ M. W. Krasny,⁸⁰ A. Krasznahorkay,³⁰
 J. K. Kraus,²¹ A. Kravchenko,²⁵ S. Kreiss,¹¹⁰ M. Kretz,^{58c} J. Kretzschmar,⁷⁴ K. Kreuzfeldt,⁵² P. Krieger,¹⁵⁸ K. Krizka,³¹
 K. Kroeninger,⁴³ H. Kroha,¹⁰¹ J. Kroll,¹²² J. Kroseberg,²¹ J. Krstic,¹³ U. Kruchonak,⁶⁵ H. Krüger,²¹ N. Krumnack,⁶⁴
 Z. V. Krumshteyn,⁶⁵ A. Kruse,¹⁷³ M. C. Kruse,⁴⁵ M. Kruskal,²² T. Kubota,⁸⁸ H. Kucuk,⁷⁸ S. Kudah,^{4b} S. Kuehn,⁴⁸
 A. Kugel,^{58c} F. Kuger,¹⁷⁴ A. Kuhl,¹³⁷ T. Kuhl,⁴² V. Kukhtin,⁶⁵ Y. Kulchitsky,⁹² S. Kuleshov,^{32b} M. Kuna,^{132a,132b} T. Kunigo,⁶⁸
 A. Kupco,¹²⁷ H. Kurashige,⁶⁷ Y. A. Kurochkin,⁹² R. Kurumida,⁶⁷ V. Kus,¹²⁷ E. S. Kuwertz,¹⁶⁹ M. Kuze,¹⁵⁷ J. Kvita,¹¹⁵
 T. Kwan,¹⁶⁹ D. Kyriazopoulos,¹³⁹ A. La Rosa,⁴⁹ J. L. La Rosa Navarro,^{24d} L. La Rotonda,^{37a,37b} C. Lacasta,¹⁶⁷
 F. Lacava,^{132a,132b} J. Lacey,²⁹ H. Lacker,¹⁶ D. Lacour,⁸⁰ V. R. Lacuesta,¹⁶⁷ E. Ladygin,⁶⁵ R. Lafaye,⁵ B. Laforge,⁸⁰
 T. Lagouri,¹⁷⁶ S. Lai,⁴⁸ L. Lambourne,⁷⁸ S. Lammers,⁶¹ C. L. Lampen,⁷ W. Lampl,⁷ E. Lançon,¹³⁶ U. Landgraf,⁴⁸
 M. P. J. Landon,⁷⁶ V. S. Lang,^{58a} J. C. Lange,¹² A. J. Lankford,¹⁶³ F. Lanni,²⁵ K. Lantzsch,³⁰ S. Laplace,⁸⁰ C. Lapoire,³⁰
 J. F. Laporte,¹³⁶ T. Lari,^{91a} F. Lasagni Manghi,^{20a,20b} M. Lassnig,³⁰ P. Laurelli,⁴⁷ W. Lavrijsen,¹⁵ A. T. Law,¹³⁷ P. Laycock,⁷⁴
 O. Le Dortz,⁸⁰ E. Le Guirriec,⁸⁵ E. Le Menedeu,¹² M. LeBlanc,¹⁶⁹ T. LeCompte,⁶ F. Ledroit-Guillon,⁵⁵ C. A. Lee,^{145b}
 S. C. Lee,¹⁵¹ L. Lee,¹ G. Lefebvre,⁸⁰ M. Lefebvre,¹⁶⁹ F. Legger,¹⁰⁰ C. Leggett,¹⁵ A. Lehan,⁷⁴ G. Lehmann Miotto,³⁰ X. Lei,⁷
 W. A. Leight,²⁹ A. Leisos,¹⁵⁴ A. G. Leister,¹⁷⁶ M. A. L. Leite,^{24d} R. Leitner,¹²⁹ D. Lellouch,¹⁷² B. Lemmer,⁵⁴
 K. J. C. Leney,⁷⁸ T. Lenz,²¹ B. Lenzi,³⁰ R. Leone,⁷ S. Leone,^{124a,124b} C. Leonidopoulos,⁴⁶ S. Leontsinis,¹⁰ C. Leroy,⁹⁵
 C. G. Lester,²⁸ M. Levchenko,¹²³ J. Levêque,⁵ D. Levin,⁸⁹ L. J. Levinson,¹⁷² M. Levy,¹⁸ A. Lewis,¹²⁰ A. M. Leyko,²¹
 M. Leyton,⁴¹ B. Li,^{33b,x} H. Li,¹⁴⁸ H. L. Li,³¹ L. Li,⁴⁵ L. Li,^{33e} S. Li,⁴⁵ Y. Li,^{33c,y} Z. Liang,¹³⁷ H. Liao,³⁴ B. Liberti,^{133a}
 A. Liblong,¹⁵⁸ P. Lichard,³⁰ K. Lie,¹⁶⁵ J. Liebal,²¹ W. Liebig,¹⁴ C. Limbach,²¹ A. Limosani,¹⁵⁰ S. C. Lin,^{151,z} T. H. Lin,⁸³
 F. Linde,¹⁰⁷ B. E. Lindquist,¹⁴⁸ J. T. Linnemann,⁹⁰ E. Lipeles,¹²² A. Lipniacka,¹⁴ M. Lisovyi,^{58b} T. M. Liss,¹⁶⁵ D. Lissauer,²⁵

A. Lister,¹⁶⁸ A. M. Litke,¹³⁷ B. Liu,^{151,aa} D. Liu,¹⁵¹ J. Liu,⁸⁵ J. B. Liu,^{33b} K. Liu,⁸⁵ L. Liu,¹⁶⁵ M. Liu,⁴⁵ M. Liu,^{33b} Y. Liu,^{33b} M. Livan,^{121a,121b} A. Lleres,⁵⁵ J. Llorente Merino,⁸² S. L. Lloyd,⁷⁶ F. Lo Sterzo,¹⁵¹ E. Lobodzinska,⁴² P. Loch,⁷ W. S. Lockman,¹³⁷ F. K. Loebinger,⁸⁴ A. E. Loevschall-Jensen,³⁶ A. Loginov,¹⁷⁶ T. Lohse,¹⁶ K. Lohwasser,⁴² M. Lokajicek,¹²⁷ B. A. Long,²² J. D. Long,⁸⁹ R. E. Long,⁷² K. A. Looper,¹¹¹ L. Lopes,^{126a} D. Lopez Mateos,⁵⁷ B. Lopez Paredes,¹³⁹ I. Lopez Paz,¹² J. Lorenz,¹⁰⁰ N. Lorenzo Martinez,⁶¹ M. Losada,¹⁶² P. Loscutoff,¹⁵ P. J. Lösel,¹⁰⁰ X. Lou,^{33a} A. Lounis,¹¹⁷ J. Love,⁶ P. A. Love,⁷² N. Lu,⁸⁹ H. J. Lubatti,¹³⁸ C. Luci,^{132a,132b} A. Lucotte,⁵⁵ F. Luehring,⁶¹ W. Lukas,⁶² L. Luminari,^{132a} O. Lundberg,^{146a,146b} B. Lund-Jensen,¹⁴⁷ D. Lynn,²⁵ R. Lysak,¹²⁷ E. Lytken,⁸¹ H. Ma,²⁵ L. L. Ma,^{33d} G. Maccarrone,⁴⁷ A. Macchiolo,¹⁰¹ C. M. Macdonald,¹³⁹ J. Machado Miguens,^{122,126b} D. Macina,³⁰ D. Madaffari,⁸⁵ R. Madar,³⁴ H. J. Maddocks,⁷² W. F. Mader,⁴⁴ A. Madsen,¹⁶⁶ S. Maeland,¹⁴ T. Maeno,²⁵ A. Maevskiy,⁹⁹ E. Magradze,⁵⁴ K. Mahboubi,⁴⁸ J. Mahlstedt,¹⁰⁷ C. Maiani,¹³⁶ C. Maidantchik,^{24a} A. A. Maier,¹⁰¹ T. Maier,¹⁰⁰ A. Maio,^{126a,126b,126d} S. Majewski,¹¹⁶ Y. Makida,⁶⁶ N. Makovec,¹¹⁷ B. Malaescu,⁸⁰ Pa. Malecki,³⁹ V. P. Maleev,¹²³ F. Malek,⁵⁵ U. Mallik,⁶³ D. Malon,⁶ C. Malone,¹⁴³ S. Maltezos,¹⁰ V. M. Malyshev,¹⁰⁹ S. Malyukov,³⁰ J. Mamuzic,⁴² G. Mancini,⁴⁷ B. Mandelli,³⁰ L. Mandelli,^{91a} I. Mandić,⁷⁵ R. Mandrysch,⁶³ J. Maneira,^{126a,126b} A. Manfredini,¹⁰¹ L. Manhaes de Andrade Filho,^{24b} J. Manjarres Ramos,^{159b} A. Mann,¹⁰⁰ P. M. Manning,¹³⁷ A. Manousakis-Katsikakis,⁹ B. Mansoulie,¹³⁶ R. Mantifel,⁸⁷ M. Mantoani,⁵⁴ L. Mapelli,³⁰ L. March,^{145c} G. Marchiori,⁸⁰ M. Marcisovsky,¹²⁷ C. P. Marino,¹⁶⁹ M. Marjanovic,¹³ F. Marroquim,^{24a} S. P. Marsden,⁸⁴ Z. Marshall,¹⁵ L. F. Marti,¹⁷ S. Marti-Garcia,¹⁶⁷ B. Martin,⁹⁰ T. A. Martin,¹⁷⁰ V. J. Martin,⁴⁶ B. Martin dit Latour,¹⁴ M. Martinez,^{12,p} S. Martin-Haugh,¹³¹ V. S. Martoiu,^{26a} A. C. Martyniuk,⁷⁸ M. Marx,¹³⁸ F. Marzano,^{132a} A. Marzin,³⁰ L. Masetti,⁸³ T. Mashimo,¹⁵⁵ R. Mashinistov,⁹⁶ J. Masik,⁸⁴ A. L. Maslennikov,^{109,d} I. Massa,^{20a,20b} L. Massa,^{20a,20b} N. Massol,⁵ P. Mastrandrea,¹⁴⁸ A. Mastroberardino,^{37a,37b} T. Masubuchi,¹⁵⁵ P. Mättig,¹⁷⁵ J. Mattmann,⁸³ J. Maurer,^{26a} S. J. Maxfield,⁷⁴ D. A. Maximov,^{109,d} R. Mazini,¹⁵¹ S. M. Mazza,^{91a,91b} L. Mazzaferro,^{133a,133b} G. Mc Goldrick,¹⁵⁸ S. P. Mc Kee,⁸⁹ A. McCarn,⁸⁹ R. L. McCarthy,¹⁴⁸ T. G. McCarthy,²⁹ N. A. McCubbin,¹³¹ K. W. McFarlane,^{56,a} J. A. Mcfayden,⁷⁸ G. Mchedlidze,⁵⁴ S. J. McMahon,¹³¹ R. A. McPherson,^{169,l} M. Medinnis,⁴² S. Meehan,^{145a} S. Mehlhase,¹⁰⁰ A. Mehta,⁷⁴ K. Meier,^{58a} C. Meineck,¹⁰⁰ B. Meirose,⁴¹ B. R. Mellado Garcia,^{145c} F. Meloni,¹⁷ A. Mengarelli,^{20a,20b} S. Menke,¹⁰¹ E. Meoni,¹⁶¹ K. M. Mercurio,⁵⁷ S. Mergelmeyer,²¹ P. Mermod,⁴⁹ L. Merola,^{104a,104b} C. Meroni,^{91a} F. S. Merritt,³¹ A. Messina,^{132a,132b} J. Metcalfe,²⁵ A. S. Mete,¹⁶³ C. Meyer,⁸³ C. Meyer,¹²² J.-P. Meyer,¹³⁶ J. Meyer,¹⁰⁷ R. P. Middleton,¹³¹ S. Miglioranza,^{164a,164c} L. Mijović,²¹ G. Mikenberg,¹⁷² M. Mikestikova,¹²⁷ M. Mikuž,⁷⁵ M. Milesi,⁸⁸ A. Milic,³⁰ D. W. Miller,³¹ C. Mills,⁴⁶ A. Milov,¹⁷² D. A. Milstead,^{146a,146b} A. A. Minaenko,¹³⁰ Y. Minami,¹⁵⁵ I. A. Minashvili,⁶⁵ A. I. Mincer,¹¹⁰ B. Mindur,^{38a} M. Mineev,⁶⁵ Y. Ming,¹⁷³ L. M. Mir,¹² T. Mitani,¹⁷¹ J. Mitrevski,¹⁰⁰ V. A. Mitsou,¹⁶⁷ A. Miucci,⁴⁹ P. S. Miyagawa,¹³⁹ J. U. Mjörnmark,⁸¹ T. Moa,^{146a,146b} K. Mochizuki,⁸⁵ S. Mohapatra,³⁵ W. Mohr,⁴⁸ S. Molander,^{146a,146b} R. Moles-Valls,¹⁶⁷ K. Mönig,⁴² C. Monini,⁵⁵ J. Monk,³⁶ E. Monnier,⁸⁵ J. Montejo Berlingen,¹² F. Monticelli,⁷¹ S. Monzani,^{132a,132b} R. W. Moore,³ N. Morange,¹¹⁷ D. Moreno,¹⁶² M. Moreno Llácer,⁵⁴ P. Morettini,^{50a} M. Morgenstern,⁴⁴ M. Morii,⁵⁷ M. Morinaga,¹⁵⁵ V. Morisbak,¹¹⁹ S. Moritz,⁸³ A. K. Morley,¹⁴⁷ G. Mornacchi,³⁰ J. D. Morris,⁷⁶ S. S. Mortensen,³⁶ A. Morton,⁵³ L. Morvaj,¹⁰³ H. G. Moser,¹⁰¹ M. Mosidze,^{51b} J. Moss,¹¹¹ K. Motohashi,¹⁵⁷ R. Mount,¹⁴³ E. Mountricha,²⁵ S. V. Mouraviev,^{96,a} E. J. W. Moyses,⁸⁶ S. Muanza,⁸⁵ R. D. Mudd,¹⁸ F. Mueller,¹⁰¹ J. Mueller,¹²⁵ K. Mueller,²¹ R. S. P. Mueller,¹⁰⁰ T. Mueller,²⁸ D. Muenstermann,⁴⁹ P. Mullen,⁵³ Y. Munwes,¹⁵³ J. A. Murillo Quijada,¹⁸ W. J. Murray,^{170,131} H. Musheghyan,⁵⁴ E. Musto,¹⁵² A. G. Myagkov,^{130,bb} M. Myska,¹²⁸ O. Nackenhorst,⁵⁴ J. Nadal,⁵⁴ K. Nagai,¹²⁰ R. Nagai,¹⁵⁷ Y. Nagai,⁸⁵ K. Nagano,⁶⁶ A. Nagarkar,¹¹¹ Y. Nagasaka,⁵⁹ K. Nagata,¹⁶⁰ M. Nagel,¹⁰¹ E. Nagy,⁸⁵ A. M. Nairz,³⁰ Y. Nakahama,³⁰ K. Nakamura,⁶⁶ T. Nakamura,¹⁵⁵ I. Nakano,¹¹² H. Namasivayam,⁴¹ R. F. Naranjo Garcia,⁴² R. Narayan,³¹ T. Naumann,⁴² G. Navarro,¹⁶² R. Nayyar,⁷ H. A. Neal,⁸⁹ P. Yu. Nechaeva,⁹⁶ T. J. Neep,⁸⁴ P. D. Nef,¹⁴³ A. Negri,^{121a,121b} M. Negrini,^{20a} S. Nektarijevic,¹⁰⁶ C. Nellist,¹¹⁷ A. Nelson,¹⁶³ S. Nemecek,¹²⁷ P. Nemethy,¹¹⁰ A. A. Nepomuceno,^{24a} M. Nessi,^{30,cc} M. S. Neubauer,¹⁶⁵ M. Neumann,¹⁷⁵ R. M. Neves,¹¹⁰ P. Nevski,²⁵ P. R. Newman,¹⁸ D. H. Nguyen,⁶ R. B. Nickerson,¹²⁰ R. Nicolaidou,¹³⁶ B. Nicquevert,³⁰ J. Nielsen,¹³⁷ N. Nikiforou,³⁵ A. Nikiforov,¹⁶ V. Nikolaenko,^{130,bb} I. Nikolic-Audit,⁸⁰ K. Nikolopoulos,¹⁸ J. K. Nilsen,¹¹⁹ P. Nilsson,²⁵ Y. Ninomiya,¹⁵⁵ A. Nisati,^{132a} R. Nisius,¹⁰¹ T. Nobe,¹⁵⁷ M. Nomachi,¹¹⁸ I. Nomidis,²⁹ T. Nooney,⁷⁶ S. Norberg,¹¹³ M. Nordberg,³⁰ O. Novgorodova,⁴⁴ S. Nowak,¹⁰¹ M. Nozaki,⁶⁶ L. Nozka,¹¹⁵ K. Ntekas,¹⁰ G. Nunes Hanninger,⁸⁸ T. Nunnemann,¹⁰⁰ E. Nurse,⁷⁸ F. Nuti,⁸⁸ B. J. O'Brien,⁴⁶ F. O'grady,⁷ D. C. O'Neil,¹⁴² V. O'Shea,⁵³ F. G. Oakham,^{29,e} H. Oberlack,¹⁰¹ T. Obermann,²¹ J. Ocariz,⁸⁰ A. Ochi,⁶⁷ I. Ochoa,⁷⁸ S. Oda,⁷⁰ S. Odaka,⁶⁶ H. Ogren,⁶¹ A. Oh,⁸⁴ S. H. Oh,⁴⁵ C. C. Ohm,¹⁵ H. Ohman,¹⁶⁶ H. Oide,³⁰ W. Okamura,¹¹⁸ H. Okawa,¹⁶⁰ Y. Okumura,³¹ T. Okuyama,¹⁵⁵ A. Olariu,^{26a} S. A. Olivares Pino,⁴⁶ D. Oliveira Damazio,²⁵ E. Oliver Garcia,¹⁶⁷ A. Olszewski,³⁹ J. Olszowska,³⁹

A. Onofre,^{126a,126e} P. U. E. Onyisi,^{31,r} C. J. Oram,^{159a} M. J. Oreglia,³¹ Y. Oren,¹⁵³ D. Orestano,^{134a,134b} N. Orlando,¹⁵⁴ C. Oropeza Barrera,⁵³ R. S. Orr,¹⁵⁸ B. Osculati,^{50a,50b} R. Ospanov,⁸⁴ G. Otero y Garzon,²⁷ H. Otono,⁷⁰ M. Ouchrif,^{135d} E. A. Ouellette,¹⁶⁹ F. Ould-Saada,¹¹⁹ A. Ouraou,¹³⁶ K. P. Oussoren,¹⁰⁷ Q. Ouyang,^{33a} A. Ovcharova,¹⁵ M. Owen,⁵³ R. E. Owen,¹⁸ V. E. Ozcan,^{19a} N. Ozturk,⁸ K. Pachal,¹⁴² A. Pacheco Pages,¹² C. Padilla Aranda,¹² M. Pagáčová,⁴⁸ S. Pagan Griso,¹⁵ E. Paganis,¹³⁹ C. Pahl,¹⁰¹ F. Paige,²⁵ P. Pais,⁸⁶ K. Pajchel,¹¹⁹ G. Palacino,^{159b} S. Palestini,³⁰ M. Palka,^{38b} D. Pallin,³⁴ A. Palma,^{126a,126b} Y. B. Pan,¹⁷³ E. Panagiotopoulou,¹⁰ C. E. Pandini,⁸⁰ J. G. Panduro Vazquez,⁷⁷ P. Pani,^{146a,146b} S. Panitkin,²⁵ L. Paolozzi,⁴⁹ Th. D. Papadopoulos,¹⁰ K. Papageorgiou,¹⁵⁴ A. Paramonov,⁶ D. Paredes Hernandez,¹⁵⁴ M. A. Parker,²⁸ K. A. Parker,¹³⁹ F. Parodi,^{50a,50b} J. A. Parsons,³⁵ U. Parzefall,⁴⁸ E. Pasqualucci,^{132a} S. Passaggio,^{50a} F. Pastore,^{134a,134b,a} Fr. Pastore,⁷⁷ G. Pásztor,²⁹ S. Patariaia,¹⁷⁵ N. D. Patel,¹⁵⁰ J. R. Pater,⁸⁴ T. Pauly,³⁰ J. Pearce,¹⁶⁹ B. Pearson,¹¹³ L. E. Pedersen,³⁶ M. Pedersen,¹¹⁹ S. Pedraza Lopez,¹⁶⁷ R. Pedro,^{126a,126b} S. V. Peleganchuk,¹⁰⁹ D. Pelikan,¹⁶⁶ H. Peng,^{33b} B. Penning,³¹ J. Penwell,⁶¹ D. V. Perepelitsa,²⁵ E. Perez Codina,^{159a} M. T. Pérez García-Estañ,¹⁶⁷ L. Perini,^{91a,91b} H. Pernegger,³⁰ S. Perrella,^{104a,104b} R. Peschke,⁴² V. D. Peshekhonov,⁶⁵ K. Peters,³⁰ R. F. Y. Peters,⁸⁴ B. A. Petersen,³⁰ T. C. Petersen,³⁶ E. Petit,⁴² A. Petridis,^{146a,146b} C. Petridou,¹⁵⁴ E. Petrolo,^{132a} F. Petrucci,^{134a,134b} N. E. Pettersson,¹⁵⁷ R. Pezoa,^{32b} P. W. Phillips,¹³¹ G. Piacquadio,¹⁴³ E. Pianori,¹⁷⁰ A. Picazio,⁴⁹ E. Piccaro,⁷⁶ M. Piccinini,^{20a,20b} M. A. Pickering,¹²⁰ R. Piegai,²⁷ D. T. Pignotti,¹¹¹ J. E. Pilcher,³¹ A. D. Pilkington,⁸⁴ J. Pina,^{126a,126b,126d} M. Pinamonti,^{164a,164c,dd} J. L. Pinfold,³ A. Pingel,³⁶ B. Pinto,^{126a} S. Pires,⁸⁰ M. Pitt,¹⁷² C. Pizio,^{91a,91b} L. Plazak,^{144a} M.-A. Pleier,²⁵ V. Pleskot,¹²⁹ E. Plotnikova,⁶⁵ P. Plucinski,^{146a,146b} D. Pluth,⁶⁴ R. Poettgen,⁸³ L. Poggioli,¹¹⁷ D. Pohl,²¹ G. Polesello,^{121a} A. Policicchio,^{37a,37b} R. Polifka,¹⁵⁸ A. Polini,^{20a} C. S. Pollard,⁵³ V. Polychronakos,²⁵ K. Pommès,³⁰ L. Pontecorvo,^{132a} B. G. Pope,⁹⁰ G. A. Popeneciu,^{26b} D. S. Popovic,¹³ A. Poppleton,³⁰ S. Pospisil,¹²⁸ K. Potamianos,¹⁵ I. N. Potrap,⁶⁵ C. J. Potter,¹⁴⁹ C. T. Potter,¹¹⁶ G. Poulard,³⁰ J. Poveda,³⁰ V. Pozdnyakov,⁶⁵ P. Pralavorio,⁸⁵ A. Pranko,¹⁵ S. Prasad,³⁰ S. Prell,⁶⁴ D. Price,⁸⁴ L. E. Price,⁶ M. Primavera,^{73a} S. Prince,⁸⁷ M. Proissl,⁴⁶ K. Prokofiev,^{60c} F. Prokoshin,^{32b} E. Protopapadaki,¹³⁶ S. Protopopescu,²⁵ J. Proudfoot,⁶ M. Przybycien,^{38a} E. Ptacek,¹¹⁶ D. Puddu,^{134a,134b} E. Pueschel,⁸⁶ D. Puldon,¹⁴⁸ M. Purohit,^{25,ee} P. Puzo,¹¹⁷ J. Qian,⁸⁹ G. Qin,⁵³ Y. Qin,⁸⁴ A. Quadt,⁵⁴ D. R. Quarrie,¹⁵ W. B. Quayle,^{164a,164b} M. Queitsch-Maitland,⁸⁴ D. Quilty,⁵³ S. Raddum,¹¹⁹ V. Radeka,²⁵ V. Radescu,⁴² S. K. Radhakrishnan,¹⁴⁸ P. Radloff,¹¹⁶ P. Rados,⁸⁸ F. Ragusa,^{91a,91b} G. Rahal,¹⁷⁸ S. Rajagopalan,²⁵ M. Rammensee,³⁰ C. Rangel-Smith,¹⁶⁶ F. Rauscher,¹⁰⁰ S. Rave,⁸³ T. Ravenscroft,⁵³ M. Raymond,³⁰ A. L. Read,¹¹⁹ N. P. Readioff,⁷⁴ D. M. Rebutti,^{121a,121b} A. Redelbach,¹⁷⁴ G. Redlinger,²⁵ R. Reece,¹³⁷ K. Reeves,⁴¹ L. Rehnisch,¹⁶ H. Reisin,²⁷ M. Relich,¹⁶³ C. Rembser,³⁰ H. Ren,^{33a} A. Renaud,¹¹⁷ M. Rescigno,^{132a} S. Resconi,^{91a} O. L. Rezanova,^{109,d} P. Reznicek,¹²⁹ R. Rezvani,⁹⁵ R. Richter,¹⁰¹ S. Richter,⁷⁸ E. Richter-Was,^{38b} O. Ricken,²¹ M. Ridel,⁸⁰ P. Rieck,¹⁶ C. J. Riegel,¹⁷⁵ J. Rieger,⁵⁴ M. Rijssenbeek,¹⁴⁸ A. Rimoldi,^{121a,121b} L. Rinaldi,^{20a} B. Ristić,⁴⁹ E. Ritsch,⁶² I. Riu,¹² F. Rizatdinova,¹¹⁴ E. Rizvi,⁷⁶ S. H. Robertson,^{87,1} A. Robichaud-Veronneau,⁸⁷ D. Robinson,²⁸ J. E. M. Robinson,⁸⁴ A. Robson,⁵³ C. Roda,^{124a,124b} S. Roe,³⁰ O. Røhne,¹¹⁹ S. Rolli,¹⁶¹ A. Romaniouk,⁹⁸ M. Romano,^{20a,20b} S. M. Romano Saez,³⁴ E. Romero Adam,¹⁶⁷ N. Rompotis,¹³⁸ M. Ronzani,⁴⁸ L. Roos,⁸⁰ E. Ros,¹⁶⁷ S. Rosati,^{132a} K. Rosbach,⁴⁸ P. Rose,¹³⁷ P. L. Rosendahl,¹⁴ O. Rosenthal,¹⁴¹ V. Rossetti,^{146a,146b} E. Rossi,^{104a,104b} L. P. Rossi,^{50a} R. Rosten,¹³⁸ M. Rotaru,^{26a} I. Roth,¹⁷² J. Rothberg,¹³⁸ D. Rousseau,¹¹⁷ C. R. Royon,¹³⁶ A. Rozanov,⁸⁵ Y. Rozen,¹⁵² X. Ruan,^{145c} F. Rubbo,¹⁴³ I. Rubinskiy,⁴² V. I. Rud,⁹⁹ C. Rudolph,⁴⁴ M. S. Rudolph,¹⁵⁸ F. Rühr,⁴⁸ A. Ruiz-Martinez,³⁰ Z. Rurikova,⁴⁸ N. A. Rusakovich,⁶⁵ A. Ruschke,¹⁰⁰ H. L. Russell,¹³⁸ J. P. Rutherford,⁷ N. Ruthmann,⁴⁸ Y. F. Ryabov,¹²³ M. Rybar,¹²⁹ G. Rybkin,¹¹⁷ N. C. Ryder,¹²⁰ A. F. Saavedra,¹⁵⁰ G. Sabato,¹⁰⁷ S. Sacerdoti,²⁷ A. Saddique,³ H. F.-W. Sadrozinski,¹³⁷ R. Sadykov,⁶⁵ F. Safai Tehrani,^{132a} M. Saimpert,¹³⁶ H. Sakamoto,¹⁵⁵ Y. Sakurai,¹⁷¹ G. Salamanna,^{134a,134b} A. Salamon,^{133a} M. Saleem,¹¹³ D. Salek,¹⁰⁷ P. H. Sales De Bruin,¹³⁸ D. Salihagic,¹⁰¹ A. Salmikov,¹⁴³ J. Salt,¹⁶⁷ D. Salvatore,^{37a,37b} F. Salvatore,¹⁴⁹ A. Salvucci,¹⁰⁶ A. Salzburger,³⁰ D. Sampsonidis,¹⁵⁴ A. Sanchez,^{104a,104b} J. Sánchez,¹⁶⁷ V. Sanchez Martinez,¹⁶⁷ H. Sandaker,¹⁴ R. L. Sandbach,⁷⁶ H. G. Sander,⁸³ M. P. Sanders,¹⁰⁰ M. Sandhoff,¹⁷⁵ C. Sandoval,¹⁶² R. Sandstroem,¹⁰¹ D. P. C. Sankey,¹³¹ M. Sannino,^{50a,50b} A. Sansoni,⁴⁷ C. Santoni,³⁴ R. Santonico,^{133a,133b} H. Santos,^{126a} I. Santoyo Castillo,¹⁴⁹ K. Sapp,¹²⁵ A. Sapronov,⁶⁵ J. G. Saraiva,^{126a,126d} B. Sarrazin,²¹ O. Sasaki,⁶⁶ Y. Sasaki,¹⁵⁵ K. Sato,¹⁶⁰ G. Sauvage,^{5,a} E. Sauvan,⁵ G. Savage,⁷⁷ P. Savard,^{158,e} C. Sawyer,¹²⁰ L. Sawyer,^{79,o} J. Saxon,³¹ C. Sbarra,^{20a} A. Sbrizzi,^{20a,20b} T. Scanlon,⁷⁸ D. A. Scannicchio,¹⁶³ M. Scarcella,¹⁵⁰ V. Scarfone,^{37a,37b} J. Schaarschmidt,¹⁷² P. Schacht,¹⁰¹ D. Schaefer,³⁰ R. Schaefer,⁴² J. Schaeffer,⁸³ S. Schaepe,²¹ S. Schaezel,^{58b} U. Schäfer,⁸³ A. C. Schaffer,¹¹⁷ D. Schaile,¹⁰⁰ R. D. Schamberger,¹⁴⁸ V. Scharf,^{58a} V. A. Schegelsky,¹²³ D. Scheirich,¹²⁹ M. Schernau,¹⁶³ C. Schiavi,^{50a,50b} C. Schillo,⁴⁸ M. Schioppa,^{37a,37b} S. Schlenker,³⁰ E. Schmidt,⁴⁸ K. Schmieden,³⁰ C. Schmitt,⁸³ S. Schmitt,^{58b} S. Schmitt,⁴² B. Schneider,^{159a} Y. J. Schnellbach,⁷⁴ U. Schnoor,⁴⁴ L. Schoeffel,¹³⁶ A. Schoening,^{58b} B. D. Schoenrock,⁹⁰ E. Schopf,²¹

A. L. S. Schorlemmer,⁵⁴ M. Schott,⁸³ D. Schouten,^{159a} J. Schovancova,⁸ S. Schramm,¹⁵⁸ M. Schreyer,¹⁷⁴ C. Schroeder,⁸³ N. Schuh,⁸³ M. J. Schultens,²¹ H.-C. Schultz-Coulon,^{58a} H. Schulz,¹⁶ M. Schumacher,⁴⁸ B. A. Schumm,¹³⁷ Ph. Schune,¹³⁶ C. Schwanenberger,⁸⁴ A. Schwartzman,¹⁴³ T. A. Schwarz,⁸⁹ Ph. Schwegler,¹⁰¹ Ph. Schwemling,¹³⁶ R. Schwienhorst,⁹⁰ J. Schwindling,¹³⁶ T. Schwindt,²¹ M. Schwoerer,⁵ F. G. Sciacca,¹⁷ E. Scifo,¹¹⁷ G. Sciolla,²³ F. Scuri,^{124a,124b} F. Scutti,²¹ J. Searcy,⁸⁹ G. Sedov,⁴² E. Sedykh,¹²³ P. Seema,²¹ S. C. Seidel,¹⁰⁵ A. Seiden,¹³⁷ F. Seifert,¹²⁸ J. M. Seixas,^{24a} G. Sekhniaidze,^{104a} K. Sekhon,⁸⁹ S. J. Sekula,⁴⁰ K. E. Selbach,⁴⁶ D. M. Seliverstov,^{123a} N. Semprini-Cesari,^{20a,20b} C. Serfon,³⁰ L. Serin,¹¹⁷ L. Serkin,^{164a,164b} T. Serre,⁸⁵ M. Sessa,^{134a,134b} R. Seuster,^{159a} H. Severini,¹¹³ T. Sfiligoj,⁷⁵ F. Sforza,¹⁰¹ A. Sfyrla,³⁰ E. Shabalina,⁵⁴ M. Shamim,¹¹⁶ L. Y. Shan,^{33a} R. Shang,¹⁶⁵ J. T. Shank,²² M. Shapiro,¹⁵ P. B. Shatalov,⁹⁷ K. Shaw,^{164a,164b} S. M. Shaw,⁸⁴ A. Shcherbakova,^{146a,146b} C. Y. Shehu,¹⁴⁹ P. Sherwood,⁷⁸ L. Shi,^{151,ff} S. Shimizu,⁶⁷ C. O. Shimmin,¹⁶³ M. Shimojima,¹⁰² M. Shiyakova,⁶⁵ A. Shmeleva,⁹⁶ D. Shoaleh Saadi,⁹⁵ M. J. Shochet,³¹ S. Shojaii,^{91a,91b} S. Shrestha,¹¹¹ E. Shulga,⁹⁸ M. A. Shupe,⁷ S. Shushkevich,⁴² P. Sicho,¹²⁷ O. Sidiropoulou,¹⁷⁴ D. Sidorov,¹¹⁴ A. Sidoti,^{20a,20b} F. Siegert,⁴⁴ Dj. Sijacki,¹³ J. Silva,^{126a,126d} Y. Silver,¹⁵³ S. B. Silverstein,^{146a} V. Simak,¹²⁸ O. Simard,⁵ Lj. Simic,¹³ S. Simion,¹¹⁷ E. Simioni,⁸³ B. Simmons,⁷⁸ D. Simon,³⁴ R. Simoniello,^{91a,91b} P. Sinervo,¹⁵⁸ N. B. Sinev,¹¹⁶ G. Siragusa,¹⁷⁴ A. N. Sisakyan,^{65a} S. Yu. Sivoklovok,⁹⁹ J. Sjölin,^{146a,146b} T. B. Sjursen,¹⁴ M. B. Skinner,⁷² H. P. Skottowe,⁵⁷ P. Skubic,¹¹³ M. Slater,¹⁸ T. Slavicek,¹²⁸ M. Slawinska,¹⁰⁷ K. Sliwa,¹⁶¹ V. Smakhtin,¹⁷² B. H. Smart,⁴⁶ L. Smestad,¹⁴ S. Yu. Smirnov,⁹⁸ Y. Smirnov,⁹⁸ L. N. Smirnova,^{99,gg} O. Smirnova,⁸¹ M. N. K. Smith,³⁵ M. Smizanska,⁷² K. Smolek,¹²⁸ A. A. Snesev,⁹⁶ G. Snidero,⁷⁶ S. Snyder,²⁵ R. Sobie,^{169,1} F. Socher,⁴⁴ A. Soffer,¹⁵³ D. A. Soh,^{151,ff} C. A. Solans,³⁰ M. Solar,¹²⁸ J. Solc,¹²⁸ E. Yu. Soldatov,⁹⁸ U. Soldevila,¹⁶⁷ A. A. Solodkov,¹³⁰ A. Soloshenko,⁶⁵ O. V. Solovyanov,¹³⁰ V. Solovyev,¹²³ P. Sommer,⁴⁸ H. Y. Song,^{33b} N. Soni,¹ A. Sood,¹⁵ A. Sopcza,¹²⁸ B. Sopko,¹²⁸ V. Sopko,¹²⁸ V. Sorin,¹² D. Sosa,^{58b} M. Sosebee,⁸ C. L. Sotiropoulou,^{124a,124b} R. Soualah,^{164a,164c} P. Soueid,⁹⁵ A. M. Soukharev,^{109,d} D. South,⁴² S. Spagnolo,^{73a,73b} M. Spalla,^{124a,124b} F. Spanò,⁷⁷ W. R. Spearman,⁵⁷ F. Spettel,¹⁰¹ R. Spighi,^{20a} G. Spigo,³⁰ L. A. Spiller,⁸⁸ M. Spousta,¹²⁹ T. Spreitzer,¹⁵⁸ R. D. St. Denis,^{53a} S. Staerz,⁴⁴ J. Stahlman,¹²² R. Stamen,^{58a} S. Stamm,¹⁶ E. Stanecka,³⁹ C. Stancu,^{134a} M. Stancu-Bellu,⁴² M. M. Stanitzki,⁴² S. Stapnes,¹¹⁹ E. A. Starchenko,¹³⁰ J. Stark,⁵⁵ P. Staroba,¹²⁷ P. Starovoitov,⁴² R. Staszewski,³⁹ P. Stavina,^{144a,a} P. Steinberg,²⁵ B. Stelzer,¹⁴² H. J. Stelzer,³⁰ O. Stelzer-Chilton,^{159a} H. Stenzel,⁵² S. Stern,¹⁰¹ G. A. Stewart,⁵³ J. A. Stillings,²¹ M. C. Stockton,⁸⁷ M. Stoebe,⁸⁷ G. Stoicea,^{26a} P. Stolte,⁵⁴ S. Stonjek,¹⁰¹ A. R. Stradling,⁸ A. Straessner,⁴⁴ M. E. Stramaglia,¹⁷ J. Strandberg,¹⁴⁷ S. Strandberg,^{146a,146b} A. Strandlie,¹¹⁹ E. Strauss,¹⁴³ M. Strauss,¹¹³ P. Strizenc,^{144b} R. Ströhmer,¹⁷⁴ D. M. Strom,¹¹⁶ R. Stroynowski,⁴⁰ A. Strubig,¹⁰⁶ S. A. Stucci,¹⁷ B. Stugu,¹⁴ N. A. Styles,⁴² D. Su,¹⁴³ J. Su,¹²⁵ R. Subramaniam,⁷⁹ A. Succurro,¹² Y. Sugaya,¹¹⁸ C. Suhr,¹⁰⁸ M. Suk,¹²⁸ V. V. Sulin,⁹⁶ S. Sultansoy,^{4c} T. Sumida,⁶⁸ S. Sun,⁵⁷ X. Sun,^{33a} J. E. Sundermann,⁴⁸ K. Suruliz,¹⁴⁹ G. Susinno,^{37a,37b} M. R. Sutton,¹⁴⁹ S. Suzuki,⁶⁶ Y. Suzuki,⁶⁶ M. Svatos,¹²⁷ S. Swedish,¹⁶⁸ M. Swiatlowski,¹⁴³ I. Sykora,^{144a} T. Sykora,¹²⁹ D. Ta,⁹⁰ C. Taccini,^{134a,134b} K. Tackmann,⁴² J. Taenzer,¹⁵⁸ A. Taffard,¹⁶³ R. Tafirout,^{159a} N. Taiblum,¹⁵³ H. Takai,²⁵ R. Takashima,⁶⁹ H. Takeda,⁶⁷ T. Takeshita,¹⁴⁰ Y. Takubo,⁶⁶ M. Talby,⁸⁵ A. A. Talyshv,^{109,d} J. Y. C. Tam,¹⁷⁴ K. G. Tan,⁸⁸ J. Tanaka,¹⁵⁵ R. Tanaka,¹¹⁷ S. Tanaka,⁶⁶ B. B. Tannenwald,¹¹¹ N. Tannoury,²¹ S. Tapprogge,⁸³ S. Tarem,¹⁵² F. Tarrade,²⁹ G. F. Tartarelli,^{91a} P. Tas,¹²⁹ M. Tasevsky,¹²⁷ T. Tashiro,⁶⁸ E. Tassi,^{37a,37b} A. Tavares Delgado,^{126a,126b} Y. Tayalati,^{135d} F. E. Taylor,⁹⁴ G. N. Taylor,⁸⁸ W. Taylor,^{159b} F. A. Teischinger,³⁰ M. Teixeira Dias Castanheira,⁷⁶ P. Teixeira-Dias,⁷⁷ K. K. Temming,⁴⁸ H. Ten Kate,³⁰ P. K. Teng,¹⁵¹ J. J. Teoh,¹¹⁸ F. Tepel,¹⁷⁵ S. Terada,⁶⁶ K. Terashi,¹⁵⁵ J. Terron,⁸² S. Terzo,¹⁰¹ M. Testa,⁴⁷ R. J. Teuscher,^{158,1} J. Therhaag,²¹ T. Thevenaux-Pelzer,³⁴ J. P. Thomas,¹⁸ J. Thomas-Wilsker,⁷⁷ E. N. Thompson,³⁵ P. D. Thompson,¹⁸ R. J. Thompson,⁸⁴ A. S. Thompson,⁵³ L. A. Thomsen,³⁶ E. Thomson,¹²² M. Thomson,²⁸ R. P. Thun,^{89,a} M. J. Tibbetts,¹⁵ R. E. Tisce Torres,⁸⁵ V. O. Tikhomirov,^{96,hh} Yu. A. Tikhonov,^{109,d} S. Timoshenko,⁹⁸ E. Tiouchichine,⁸⁵ P. Tipton,¹⁷⁶ S. Tisserant,⁸⁵ T. Todorov,^{5a} S. Todorova-Nova,¹²⁹ J. Tojo,⁷⁰ S. Tokár,^{144a} K. Tokushuku,⁶⁶ K. Tollefson,⁹⁰ E. Tolley,⁵⁷ L. Tomlinson,⁸⁴ M. Tomoto,¹⁰³ L. Tompkins,^{143,ii} K. Toms,¹⁰⁵ E. Torrence,¹¹⁶ H. Torres,¹⁴² E. Torró Pastor,¹⁶⁷ J. Toth,^{85,jj} F. Touchard,⁸⁵ D. R. Tovey,¹³⁹ T. Trefzger,¹⁷⁴ L. Tremblet,³⁰ A. Tricoli,³⁰ I. M. Trigger,^{159a} S. Trincaz-Duvoid,⁸⁰ M. F. Tripiana,¹² W. Trischuk,¹⁵⁸ B. Trocmé,⁵⁵ C. Troncon,^{91a} M. Trotter-McDonald,¹⁵ M. Trovatelli,^{134a,134b} P. True,⁹⁰ L. Truong,^{164a,164c} M. Trzebinski,³⁹ A. Trzupek,³⁹ C. Tsarouchas,³⁰ J. C.-L. Tseng,¹²⁰ P. V. Tsiarshka,⁹² D. Tsionou,¹⁵⁴ G. Tsipolitis,¹⁰ N. Tsirintanis,⁹ S. Tsiskaridze,¹² V. Tsiskaridze,⁴⁸ E. G. Tskhadadze,^{51a} I. I. Tsukerman,⁹⁷ V. Tsulaia,¹⁵ S. Tsuno,⁶⁶ D. Tsybychev,¹⁴⁸ A. Tudorache,^{26a} V. Tudorache,^{26a} A. N. Tuna,¹²² S. A. Tuppiti,^{20a,20b} S. Turchikhin,^{99,gg} D. Turecek,¹²⁸ R. Turra,^{91a,91b} A. J. Turvey,⁴⁰ P. M. Tuts,³⁵ A. Tykhonov,⁴⁹ M. Tylmad,^{146a,146b} M. Tyndel,¹³¹ I. Ueda,¹⁵⁵ R. Ueno,²⁹ M. Ughetto,^{146a,146b} M. Ugland,¹⁴ M. Uhlenbrock,²¹ F. Ukegawa,¹⁶⁰ G. Unal,³⁰ A. Undrus,²⁵ G. Unel,¹⁶³ F. C. Ungaro,⁴⁸ Y. Unno,⁶⁶ C. Unverdorben,¹⁰⁰ J. Urban,^{144b}

P. Urquijo,⁸⁸ P. Urrejola,⁸³ G. Usai,⁸ A. Usanova,⁶² L. Vacavant,⁸⁵ V. Vacek,¹²⁸ B. Vachon,⁸⁷ C. Valderanis,⁸³ N. Valencic,¹⁰⁷ S. Valentini,^{20a,20b} A. Valero,¹⁶⁷ L. Valery,¹² S. Valkar,¹²⁹ E. Valladolid Gallego,¹⁶⁷ S. Vallecorsa,⁴⁹ J. A. Valls Ferrer,¹⁶⁷ W. Van Den Wollenberg,¹⁰⁷ P. C. Van Der Deijl,¹⁰⁷ R. van der Geer,¹⁰⁷ H. van der Graaf,¹⁰⁷ R. Van Der Leeuw,¹⁰⁷ N. van Eldik,¹⁵² P. van Gemmeren,⁶ J. Van Nieuwkoop,¹⁴² I. van Vulpen,¹⁰⁷ M. C. van Woerden,³⁰ M. Vanadia,^{132a,132b} W. Vandelli,³⁰ R. Vanguri,¹²² A. Vaniachine,⁶ F. Vannucci,⁸⁰ G. Vardanyan,¹⁷⁷ R. Vari,^{132a} E. W. Varnes,⁷ T. Varol,⁴⁰ D. Varouchas,⁸⁰ A. Vartapetian,⁸ K. E. Varvell,¹⁵⁰ F. Vazeille,³⁴ T. Vazquez Schroeder,⁸⁷ J. Veatch,⁷ F. Veloso,^{126a,126c} T. Velz,²¹ S. Veneziano,^{132a} A. Ventura,^{73a,73b} D. Ventura,⁸⁶ M. Venturi,¹⁶⁹ N. Venturi,¹⁵⁸ A. Venturini,²³ V. Vercesi,^{121a} M. Verducci,^{132a,132b} W. Verkerke,¹⁰⁷ J. C. Vermeulen,¹⁰⁷ A. Vest,⁴⁴ M. C. Vetterli,^{142,e} O. Viazlo,⁸¹ I. Vichou,¹⁶⁵ T. Vickey,¹³⁹ O. E. Vickey Boeriu,¹³⁹ G. H. A. Viehhauser,¹²⁰ S. Viel,¹⁵ R. Vigne,³⁰ M. Villa,^{20a,20b} M. Villaplana Perez,^{91a,91b} E. Vilucchi,⁴⁷ M. G. Vinciter,²⁹ V. B. Vinogradov,⁶⁵ I. Vivarelli,¹⁴⁹ F. Vives Vaque,³ S. Vlachos,¹⁰ D. Vladoiu,¹⁰⁰ M. Vlasak,¹²⁸ M. Vogel,^{32a} P. Vokac,¹²⁸ G. Volpi,^{124a,124b} M. Volpi,⁸⁸ H. von der Schmitt,¹⁰¹ H. von Radziewski,⁴⁸ E. von Toerne,²¹ V. Vorobel,¹²⁹ K. Vorobev,⁹⁸ M. Vos,¹⁶⁷ R. Voss,³⁰ J. H. Vosseveld,⁷⁴ N. Vranjes,¹³ M. Vranjes Milosavljevic,¹³ V. Vrba,¹²⁷ M. Vreeswijk,¹⁰⁷ R. Vuillermet,³⁰ I. Vukotic,³¹ Z. Vykydal,¹²⁸ P. Wagner,²¹ W. Wagner,¹⁷⁵ H. Wahlberg,⁷¹ S. Wahrmund,⁴⁴ J. Wakabayashi,¹⁰³ J. Walder,⁷² R. Walker,¹⁰⁰ W. Walkowiak,¹⁴¹ C. Wang,^{33c} F. Wang,¹⁷³ H. Wang,¹⁵ H. Wang,⁴⁰ J. Wang,⁴² J. Wang,^{33a} K. Wang,⁸⁷ R. Wang,⁶ S. M. Wang,¹⁵¹ T. Wang,²¹ X. Wang,¹⁷⁶ C. Wanotayaroj,¹¹⁶ A. Warburton,⁸⁷ C. P. Ward,²⁸ D. R. Wardrope,⁷⁸ M. Warsinsky,⁴⁸ A. Washbrook,⁴⁶ C. Wasicki,⁴² P. M. Watkins,¹⁸ A. T. Watson,¹⁸ I. J. Watson,¹⁵⁰ M. F. Watson,¹⁸ G. Watts,¹³⁸ S. Watts,⁸⁴ B. M. Waugh,⁷⁸ S. Webb,⁸⁴ M. S. Weber,¹⁷ S. W. Weber,¹⁷⁴ J. S. Webster,³¹ A. R. Weidberg,¹²⁰ B. Weinert,⁶¹ J. Weingarten,⁵⁴ C. Weiser,⁴⁸ H. Weits,¹⁰⁷ P. S. Wells,³⁰ T. Wenaus,²⁵ T. Wengler,³⁰ S. Wenig,³⁰ N. Wermes,²¹ M. Werner,⁴⁸ P. Werner,³⁰ M. Wessels,^{58a} J. Wetter,¹⁶¹ K. Whalen,²⁹ A. M. Wharton,⁷² A. White,⁸ M. J. White,¹ R. White,^{32b} S. White,^{124a,124b} D. Whiteson,¹⁶³ F. J. Wickens,¹³¹ W. Wiedenmann,¹⁷³ M. Wielers,¹³¹ P. Wienemann,²¹ C. Wiglesworth,³⁶ L. A. M. Wiik-Fuchs,²¹ A. Wildauer,¹⁰¹ H. G. Wilkens,³⁰ H. H. Williams,¹²² S. Williams,¹⁰⁷ C. Willis,⁹⁰ S. Willocq,⁸⁶ A. Wilson,⁸⁹ J. A. Wilson,¹⁸ I. Wingerter-Seez,⁵ F. Winklmeier,¹¹⁶ B. T. Winter,²¹ M. Wittgen,¹⁴³ J. Wittkowski,¹⁰⁰ S. J. Wollstadt,⁸³ M. W. Wolter,³⁹ H. Wolters,^{126a,126c} B. K. Wosiek,³⁹ J. Wotschack,³⁰ M. J. Woudstra,⁸⁴ K. W. Wozniak,³⁹ M. Wu,⁵⁵ M. Wu,³¹ S. L. Wu,¹⁷³ X. Wu,⁴⁹ Y. Wu,⁸⁹ T. R. Wyatt,⁸⁴ B. M. Wynne,⁴⁶ S. Xella,³⁶ D. Xu,^{33a} L. Xu,^{33b,kk} B. Yabsley,¹⁵⁰ S. Yacoob,^{145b,ll} R. Yakabe,⁶⁷ M. Yamada,⁶⁶ Y. Yamaguchi,¹¹⁸ A. Yamamoto,⁶⁶ S. Yamamoto,¹⁵⁵ T. Yamanaka,¹⁵⁵ K. Yamauchi,¹⁰³ Y. Yamazaki,⁶⁷ Z. Yan,²² H. Yang,^{33e} H. Yang,¹⁷³ Y. Yang,¹⁵¹ L. Yao,^{33a} W.-M. Yao,¹⁵ Y. Yasu,⁶⁶ E. Yatsenko,⁵ K. H. Yau Wong,²¹ J. Ye,⁴⁰ S. Ye,²⁵ I. Yeletsikh,⁶⁵ A. L. Yen,⁵⁷ E. Yildirim,⁴² K. Yorita,¹⁷¹ R. Yoshida,⁶ K. Yoshihara,¹²² C. Young,¹⁴³ C. J. S. Young,³⁰ S. Youssef,²² D. R. Yu,¹⁵ J. Yu,⁸ J. M. Yu,⁸⁹ J. Yu,¹¹⁴ L. Yuan,⁶⁷ A. Yurkewicz,¹⁰⁸ I. Yusuff,^{28,mm} B. Zabinski,³⁹ R. Zaidan,⁶³ A. M. Zaitsev,^{130,bb} J. Zalieckas,¹⁴ A. Zaman,¹⁴⁸ S. Zambito,⁵⁷ L. Zanello,^{132a,132b} D. Zanzi,⁸⁸ C. Zeitnitz,¹⁷⁵ M. Zeman,¹²⁸ A. Zemla,^{38a} K. Zengel,²³ O. Zenin,¹³⁰ T. Ženiš,^{144a} D. Zerwas,¹¹⁷ D. Zhang,⁸⁹ F. Zhang,¹⁷³ J. Zhang,⁶ L. Zhang,⁴⁸ R. Zhang,^{33b} X. Zhang,^{33d} Z. Zhang,¹¹⁷ X. Zhao,⁴⁰ Y. Zhao,^{33d,117} Z. Zhao,^{33b} A. Zhemchugov,⁶⁵ J. Zhong,¹²⁰ B. Zhou,⁸⁹ C. Zhou,⁴⁵ L. Zhou,³⁵ L. Zhou,⁴⁰ N. Zhou,¹⁶³ C. G. Zhu,^{33d} H. Zhu,^{33a} J. Zhu,⁸⁹ Y. Zhu,^{33b} X. Zhuang,^{33a} K. Zhukov,⁹⁶ A. Zibell,¹⁷⁴ D. Zieminska,⁶¹ N. I. Zimine,⁶⁵ C. Zimmermann,⁸³ S. Zimmermann,⁴⁸ Z. Zinonos,⁵⁴ M. Zinser,⁸³ M. Ziolkowski,¹⁴¹ L. Živković,¹³ G. Zoernig,¹⁷³ A. Zoccoli,^{20a,20b} M. zur Nedden,¹⁶ G. Zurzolo,^{104a,104b} and L. Zwalinski³⁰

(ATLAS Collaboration)

¹*Department of Physics, University of Adelaide, Adelaide, Australia*²*Physics Department, SUNY Albany, Albany, New York, USA*³*Department of Physics, University of Alberta, Edmonton, Alberta, Canada*^{4a}*Department of Physics, Ankara University, Ankara, Turkey*^{4b}*Istanbul Aydin University, Istanbul, Turkey*^{4c}*Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey*⁵*LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France*⁶*High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA*⁷*Department of Physics, University of Arizona, Tucson, Arizona, USA*⁸*Department of Physics, The University of Texas at Arlington, Arlington, Texas, USA*⁹*Physics Department, University of Athens, Athens, Greece*¹⁰*Physics Department, National Technical University of Athens, Zografou, Greece*¹¹*Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan*

- ¹²*Institut de Física d'Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain*
- ¹³*Institute of Physics, University of Belgrade, Belgrade, Serbia*
- ¹⁴*Department for Physics and Technology, University of Bergen, Bergen, Norway*
- ¹⁵*Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, California, USA*
- ¹⁶*Department of Physics, Humboldt University, Berlin, Germany*
- ¹⁷*Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland*
- ¹⁸*School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom*
- ^{19a}*Department of Physics, Bogazici University, Istanbul, Turkey*
- ^{19b}*Department of Physics, Dogus University, Istanbul, Turkey*
- ^{19c}*Department of Physics Engineering, Gaziantep University, Gaziantep, Turkey*
- ^{20a}*INFN Sezione di Bologna, Italy*
- ^{20b}*Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy*
- ²¹*Physikalisches Institut, University of Bonn, Bonn, Germany*
- ²²*Department of Physics, Boston University, Boston, Massachusetts, USA*
- ²³*Department of Physics, Brandeis University, Waltham, Massachusetts, USA*
- ^{24a}*Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil*
- ^{24b}*Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora, Brazil*
- ^{24c}*Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei, Brazil*
- ^{24d}*Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil*
- ²⁵*Physics Department, Brookhaven National Laboratory, Upton, New York, USA*
- ^{26a}*National Institute of Physics and Nuclear Engineering, Bucharest, Romania*
- ^{26b}*National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca, Romania*
- ^{26c}*University Politehnica Bucharest, Bucharest, Romania*
- ^{26d}*West University in Timisoara, Timisoara, Romania*
- ²⁷*Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina*
- ²⁸*Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom*
- ²⁹*Department of Physics, Carleton University, Ottawa, Ontario, Canada*
- ³⁰*CERN, Geneva, Switzerland*
- ³¹*Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA*
- ^{32a}*Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile*
- ^{32b}*Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile*
- ^{33a}*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China*
- ^{33b}*Department of Modern Physics, University of Science and Technology of China, Anhui, China*
- ^{33c}*Department of Physics, Nanjing University, Jiangsu, China*
- ^{33d}*School of Physics, Shandong University, Shandong, China*
- ^{33e}*Department of Physics and Astronomy, Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University, Shanghai, China*
- ^{33f}*Physics Department, Tsinghua University, Beijing 100084, China*
- ³⁴*Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France*
- ³⁵*Nevis Laboratory, Columbia University, Irvington, New York, USA*
- ³⁶*Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark*
- ^{37a}*INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy*
- ^{37b}*Dipartimento di Fisica, Università della Calabria, Rende, Italy*
- ^{38a}*AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland*
- ^{38b}*Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland*
- ³⁹*Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland*
- ⁴⁰*Physics Department, Southern Methodist University, Dallas, Texas, USA*
- ⁴¹*Physics Department, University of Texas at Dallas, Richardson, Texas, USA*
- ⁴²*DESY, Hamburg and Zeuthen, Germany*
- ⁴³*Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany*
- ⁴⁴*Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany*
- ⁴⁵*Department of Physics, Duke University, Durham, North Carolina, USA*
- ⁴⁶*SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom*
- ⁴⁷*INFN Laboratori Nazionali di Frascati, Frascati, Italy*
- ⁴⁸*Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany*
- ⁴⁹*Section de Physique, Université de Genève, Geneva, Switzerland*
- ^{50a}*INFN Sezione di Genova, Italy*
- ^{50b}*Dipartimento di Fisica, Università di Genova, Genova, Italy*

- ^{51a}*E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia*
^{51b}*High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*
⁵²*II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany*
⁵³*SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom*
⁵⁴*II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany*
⁵⁵*Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France*
⁵⁶*Department of Physics, Hampton University, Hampton, Virginia, USA*
⁵⁷*Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA*
^{58a}*Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
^{58b}*Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
^{58c}*ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany*
⁵⁹*Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan*
^{60a}*Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China*
^{60b}*Department of Physics, The University of Hong Kong, Hong Kong, China*
^{60c}*Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China*
⁶¹*Department of Physics, Indiana University, Bloomington, Indiana, USA*
⁶²*Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria*
⁶³*University of Iowa, Iowa City, Iowa, USA*
⁶⁴*Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA*
⁶⁵*Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia*
⁶⁶*KEK, High Energy Accelerator Research Organization, Tsukuba, Japan*
⁶⁷*Graduate School of Science, Kobe University, Kobe, Japan*
⁶⁸*Faculty of Science, Kyoto University, Kyoto, Japan*
⁶⁹*Kyoto University of Education, Kyoto, Japan*
⁷⁰*Department of Physics, Kyushu University, Fukuoka, Japan*
⁷¹*Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina*
⁷²*Physics Department, Lancaster University, Lancaster, United Kingdom*
^{73a}*INFN Sezione di Lecce, Italy*
^{73b}*Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy*
⁷⁴*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*
⁷⁵*Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia*
⁷⁶*School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom*
⁷⁷*Department of Physics, Royal Holloway University of London, Surrey, United Kingdom*
⁷⁸*Department of Physics and Astronomy, University College London, London, United Kingdom*
⁷⁹*Louisiana Tech University, Ruston, Louisiana, USA*
⁸⁰*Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France*
⁸¹*Fysiska institutionen, Lunds universitet, Lund, Sweden*
⁸²*Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain*
⁸³*Institut für Physik, Universität Mainz, Mainz, Germany*
⁸⁴*School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*
⁸⁵*CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France*
⁸⁶*Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA*
⁸⁷*Department of Physics, McGill University, Montreal, Québec, Canada*
⁸⁸*School of Physics, University of Melbourne, Victoria, Australia*
⁸⁹*Department of Physics, The University of Michigan, Ann Arbor, Michigan, USA*
⁹⁰*Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA*
^{91a}*INFN Sezione di Milano, Italy*
^{91b}*Dipartimento di Fisica, Università di Milano, Milano, Italy*
⁹²*B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus*
⁹³*National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus*
⁹⁴*Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts, USA*
⁹⁵*Group of Particle Physics, University of Montreal, Montreal, Québec, Canada*
⁹⁶*P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia*
⁹⁷*Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia*
⁹⁸*National Research Nuclear University MEPhI, Moscow, Russia*
⁹⁹*D.V. Skobel'syn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia*
¹⁰⁰*Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany*
¹⁰¹*Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany*
¹⁰²*Nagasaki Institute of Applied Science, Nagasaki, Japan*
¹⁰³*Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan*

- ^{104a}*INFN Sezione di Napoli, Italy*
- ^{104b}*Dipartimento di Fisica, Università di Napoli, Napoli, Italy*
- ¹⁰⁵*Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA*
- ¹⁰⁶*Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands*
- ¹⁰⁷*Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands*
- ¹⁰⁸*Department of Physics, Northern Illinois University, DeKalb, Illinois, USA*
- ¹⁰⁹*Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia*
- ¹¹⁰*Department of Physics, New York University, New York, New York, USA*
- ¹¹¹*Ohio State University, Columbus, Ohio, USA*
- ¹¹²*Faculty of Science, Okayama University, Okayama, Japan*
- ¹¹³*Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA*
- ¹¹⁴*Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA*
- ¹¹⁵*Palacký University, RCPTM, Olomouc, Czech Republic*
- ¹¹⁶*Center for High Energy Physics, University of Oregon, Eugene, Oregon, USA*
- ¹¹⁷*LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France*
- ¹¹⁸*Graduate School of Science, Osaka University, Osaka, Japan*
- ¹¹⁹*Department of Physics, University of Oslo, Oslo, Norway*
- ¹²⁰*Department of Physics, Oxford University, Oxford, United Kingdom*
- ^{121a}*INFN Sezione di Pavia, Italy*
- ^{121b}*Dipartimento di Fisica, Università di Pavia, Pavia, Italy*
- ¹²²*Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA*
- ¹²³*Petersburg Nuclear Physics Institute, Gatchina, Russia*
- ^{124a}*INFN Sezione di Pisa, Italy*
- ^{124b}*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
- ¹²⁵*Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA*
- ^{126a}*Laboratorio de Instrumentacao e Fisica Experimental de Particulas - LIP, Lisboa, Portugal*
- ^{126b}*Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal*
- ^{126c}*Department of Physics, University of Coimbra, Coimbra, Portugal*
- ^{126d}*Centro de Fisica Nuclear da Universidade de Lisboa, Lisboa, Portugal*
- ^{126e}*Departamento de Fisica, Universidade do Minho, Braga, Portugal*
- ^{126f}*Departamento de Fisica Teorica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain*
- ^{126g}*Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal*
- ¹²⁷*Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic*
- ¹²⁸*Czech Technical University in Prague, Praha, Czech Republic*
- ¹²⁹*Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic*
- ¹³⁰*State Research Center Institute for High Energy Physics, Protvino, Russia*
- ¹³¹*Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ^{132a}*INFN Sezione di Roma, Italy*
- ^{132b}*Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy*
- ^{133a}*INFN Sezione di Roma Tor Vergata, Italy*
- ^{133b}*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
- ^{134a}*INFN Sezione di Roma Tre, Italy*
- ^{134b}*Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy*
- ^{135a}*Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca, Morocco*
- ^{135b}*Centre National de l'Energie des Sciences Techniques Nucleaires, Rabat, Morocco*
- ^{135c}*Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco*
- ^{135d}*Faculté des Sciences, Université Mohamed Premier and LTPM, Oujda, Morocco*
- ^{135e}*Faculté des sciences, Université Mohammed V-Agdal, Rabat, Morocco*
- ¹³⁶*DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat à l'Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France*
- ¹³⁷*Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, California, USA*
- ¹³⁸*Department of Physics, University of Washington, Seattle, Washington, USA*
- ¹³⁹*Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
- ¹⁴⁰*Department of Physics, Shinshu University, Nagano, Japan*
- ¹⁴¹*Fachbereich Physik, Universität Siegen, Siegen, Germany*
- ¹⁴²*Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada*
- ¹⁴³*SLAC National Accelerator Laboratory, Stanford, California, USA*
- ^{144a}*Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava, Slovak Republic*
- ^{144b}*Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*

- ^{145a}*Department of Physics, University of Cape Town, Cape Town, South Africa*
^{145b}*Department of Physics, University of Johannesburg, Johannesburg, South Africa*
^{145c}*School of Physics, University of the Witwatersrand, Johannesburg, South Africa*
^{146a}*Department of Physics, Stockholm University, Sweden*
^{146b}*The Oskar Klein Centre, Stockholm, Sweden*
¹⁴⁷*Physics Department, Royal Institute of Technology, Stockholm, Sweden*
¹⁴⁸*Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, New York, USA*
¹⁴⁹*Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
¹⁵⁰*School of Physics, University of Sydney, Sydney, Australia*
¹⁵¹*Institute of Physics, Academia Sinica, Taipei, Taiwan*
¹⁵²*Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel*
¹⁵³*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
¹⁵⁴*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
¹⁵⁵*International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan*
¹⁵⁶*Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan*
¹⁵⁷*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
¹⁵⁸*Department of Physics, University of Toronto, Toronto, Ontario, Canada*
^{159a}*TRIUMF, Vancouver, British Columbia, Canada*
^{159b}*Department of Physics and Astronomy, York University, Toronto, Ontario, Canada*
¹⁶⁰*Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan*
¹⁶¹*Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA*
¹⁶²*Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia*
¹⁶³*Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA*
^{164a}*INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy*
^{164b}*ICTP, Trieste, Italy*
^{164c}*Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy*
¹⁶⁵*Department of Physics, University of Illinois, Urbana, Illinois, USA*
¹⁶⁶*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
¹⁶⁷*Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain*
¹⁶⁸*Department of Physics, University of British Columbia, Vancouver, British Columbia, Canada*
¹⁶⁹*Department of Physics and Astronomy, University of Victoria, Victoria, British Columbia, Canada*
¹⁷⁰*Department of Physics, University of Warwick, Coventry, United Kingdom*
¹⁷¹*Waseda University, Tokyo, Japan*
¹⁷²*Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel*
¹⁷³*Department of Physics, University of Wisconsin, Madison, Wisconsin, USA*
¹⁷⁴*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany*
¹⁷⁵*Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
¹⁷⁶*Department of Physics, Yale University, New Haven, Connecticut, USA*
¹⁷⁷*Yerevan Physics Institute, Yerevan, Armenia*
¹⁷⁸*Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France*

^aDeceased.

^bAlso at Department of Physics, King's College London, London, United Kingdom.

^cAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

^dAlso at Novosibirsk State University, Novosibirsk, Russia.

^eAlso at TRIUMF, Vancouver BC, Canada.

^fAlso at Department of Physics, California State University, Fresno CA, USA.

^gAlso at Department of Physics, University of Fribourg, Fribourg, Switzerland.

^hAlso at Departamento de Física e Astronomia, Faculdade de Ciências, Universidade do Porto, Portugal.

ⁱAlso at Tomsk State University, Tomsk, Russia.

^jAlso at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.

^kAlso at Università di Napoli Parthenope, Napoli, Italy.

^lAlso at Institute of Particle Physics (IPP), Canada.

^mAlso at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.

ⁿAlso at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.

^oAlso at Louisiana Tech University, Ruston LA, USA.

^pAlso at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

^qAlso at Department of Physics, National Tsing Hua University, Taiwan.

^rAlso at Department of Physics, The University of Texas at Austin, Austin TX, USA.

- ^s Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.
- ^t Also at CERN, Geneva, Switzerland.
- ^u Also at Georgian Technical University (GTU), Tbilisi, Georgia.
- ^v Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan.
- ^w Also at Manhattan College, New York NY, USA.
- ^x Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
- ^y Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.
- ^z Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.
- ^{aa} Also at School of Physics, Shandong University, Shandong, China.
- ^{bb} Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.
- ^{cc} Also at Section de Physique, Université de Genève, Geneva, Switzerland.
- ^{dd} Also at International School for Advanced Studies (SISSA), Trieste, Italy.
- ^{ee} Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, USA.
- ^{ff} Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.
- ^{gg} Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia.
- ^{hh} Also at National Research Nuclear University MEPhI, Moscow, Russia.
- ⁱⁱ Also at Department of Physics, Stanford University, Stanford CA, USA.
- ^{jj} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
- ^{kk} Also at Department of Physics, The University of Michigan, Ann Arbor MI, USA.
- ^{ll} Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa.
- ^{mmm} Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.